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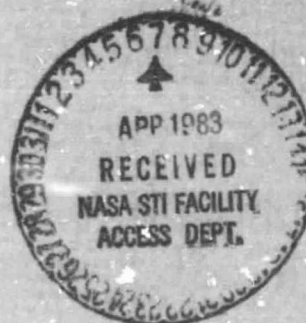
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ATMOSPHERIC RADIATION
MODEL FOR
WATER SURFACES

FINAL TECHNICAL REPORT

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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) The purpose of this investigation is to extend an atmospheric correction model to account for various atmospheric radiation components in remotely sensed data. A previous study considered the path radiance which results from singly and multiply scattered radiation along the line of sight from a multispectral aircraft sensor and a target on a water surface. This report considers		

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additional components such as the atmospheric path radiance which results from singly-scattered sky radiation specularly reflected by the water surface. It also considers a component which is referred to as the virtual sun path radiance, i.e. the singly-scattered path radiance which results from the solar radiation which is specularly reflected by the water surface.

These atmospheric radiation components are coded into a computer program for the analysis of multispectral remote sensor data over the Great Lakes of the United States. The user must know certain parameters, such as the visibility or spectral optical thickness of the atmosphere and the geometry of the sensor with respect to the sun and the target elements under investigation.

Suggestions and recommendations are given for further investigation of the problem of the remote sensing of water surfaces. If all of these extrinsic radiation components are properly accounted for, then the intrinsic water radiance can be found by applying the algorithm or an adaptation of the algorithm in this report. As a result, one would then be able to know the actual surface water spectral radiation field independent of the atmosphere.

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PREFACE

This report describes part of the work done on a research program in the remote sensing of the Great Lakes using a multi-spectral scanner aboard an aircraft. The research has been conducted for the NASA-Lewis Research Center, Cleveland, Ohio, by Science Applications, Inc., at the Dayton, Ohio, office. The primary objective of this program is to develop remote sensing as a practical technique for the analysis of the Great Lakes.

Remote sensing of the environment involves the transfer of radiation from the Earth's surface through the atmosphere to a sensor which is located at some point within the atmosphere. For water surfaces with their inherently low reflectances, the atmospheric scattering of solar radiation acts as a significant noise factor. In this report we have extended an existing model to include various atmospheric radiation components so that the resulting mathematical algorithm will allow one to extract a radiance value which is more nearly representative of the actual radiance of the water, independent of atmospheric effects.

This research was performed under contract NAS3-22495 and covers the period from 4 March 1981 through 15 July 1982. Mr. Thom Coney served as Technical Monitor of the contract and Dr. Robert E. Turner of Science Applications, Inc., was the program manager and principal investigator.

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SYMBOLS

E_0	solar irradiance at the top of the atmosphere
E_0'	solar irradiance of the virtual sun
$L_0(i)$	radiance incident at angle i
$L_s(i)$	surface radiance at angle i
L_{SKYRE}	reflected sky radiance
L_{SOLRE}	reflected solar radiance
L_{SKY}	sky radiance
L_{SOL}	direct beam radiance
L_{PMS}	multiply scattered path radiance
L_{PSS}	singly scattered path radiance
L_U	upward scattered radiance beneath water surface
L_{INIR}	radiance scattered from beneath water surface
L_{SKYSS}	singly scattered sky radiance
L_{PTOT}	total radiance
L_{VSOL}	radiance from virtual sun
L_{VP}	virtual sun path radiance
n	index of refraction of water relative to air
p	slope probability
$p(x)$	scattering phase function
T	transmittance

(GREEK SYMBOLS)

θ_i	incident angle
θ_r	reflected angle
ρ_s	Fresnel reflectance

$\rho_f(i)$	Fresnel reflectance at angle i
θ	nadir view angle
θ_0	solar zenith angle
ϕ	azimuth view angle
ϕ_0	solar azimuth angle
μ	cosine of nadir view angle
μ_0	cosine of solar zenith angle
χ_{SKY}	scattering angle
τ	optical depth
τ_0	optical thickness
η	forward scattering parameter
ρ	surface albedo
ω_0	single scattering albedo

SUMMARY

In the analysis of remotely sensed data on bodies of water, the atmosphere obscures the inherent surface features as a result of the scattering and absorption of solar radiation. In the case of multispectral data acquired by aircraft or spacecraft sensors, one can preprocess the data by applying mathematical models and algorithms to the digitized data. The mathematical model developed in this investigation is specifically designed to account for various components of the visible and infrared radiation in the atmosphere which interfere with the inherent signal from a water surface. If the atmospheric parameters are known, then when the algorithm is applied to the multispectral data sets, an improved or corrected data set will result.

This improved atmospheric correction model allows for the path radiance in the atmosphere as a result of singly-scattered solar radiation and also singly-scattered solar-reflected radiation. In addition, the model includes a singly-scattered sky radiation component for the radiation which is reflected by the water surface. Comparisons are made among the relative magnitudes of these radiation components in terms of the geometric and environmental factors. Recommendations are presented for a more advanced model which would include the corresponding radiation components for multiple scattering.

INTRODUCTION

Multispectral scanner data obtained by sensors aboard aircraft and spacecraft allow a user to examine the detailed physical properties of a surface. These properties are of interest to many investigations in various disciplines such as land use studies, agriculture, hydrology, forestry, and oceanography. In all of these investigations, however, the scattering of visible and infrared radiation by the atmospheric constituents will reduce the inherent surface radiance and add a path radiance to the attenuated radiance from the target. For many cases of the remote sensing of bright land areas on relatively clear days the attenuated radiation from the surface is rather large as compared to the atmospheric path radiance. For water bodies, with inherently low reflectances, this is no longer true and the path radiance can be a major effect in the total radiance at the sensor.

The purpose of this investigation was to extend an existing atmospheric radiative transfer model to include other radiation components which did not exist in the previous model. These additional atmospheric radiation components include specific effects for the remote sensing of water surfaces. The model is used in conjunction with an algorithm specifically designed for the analysis of multispectral data.

The determination of the atmospheric radiation components is important for the analysis of the probability of misclassification of various classes of surface materials. To first order one may consider the so called linear transfer problem in which the path radiance is constant over varying surface reflectances for a horizontally spatially uniform haze. For a non-uniform haze, however, the path radiance can vary, thereby resulting in a higher

probability of misclassification of objects if the degree of non-uniformity is unknown. A second-order effect, but one which can become quite important for the remote sensing of high-contrast targets is the adjacency effect. This is when radiation from a bright target causes an increase in the path radiance with respect to the radiance from a neighboring dark target. This problem would exist, for example, in the remote sensing of water bodies near bright sandy beaches. The results should be evident in the brightened image of the water near the shoreline, provided the effects of waves and whitecaps are eliminated. This second-order adjacency effect is not included in the model or algorithm in this investigation but the effect can be accounted for if the investigator has sufficiently detailed atmospheric data on the horizontal stratification of aerosols.

Multiple scattering is particularly important if the sky is hazy. These effects are considered in the model for path radiance which results from the sun as a source. We have not included the multiple scattering effects for solar-reflected radiation.

OPTICAL PROPERTIES OF WATER

The interpretation of remote sensing data collected over water surfaces requires a detailed knowledge of the optical properties of water and the air-water interface. Water is unusual as a natural surface because it is a specular reflector and because in the visible regions sensible data can be obtained from well below the water surface [1]. Also of importance is the phenomenon of refraction, which occurs when radiation passes through the air-water boundary. These properties of water and their significance in terms of remote sensing are discussed in detail in this chapter.

3.1 REFLECTION, REFLECTANCE AND REFRACTION

Most natural surfaces are approximately Lambertian--reflecting incident radiance equally in all directions. A smooth water surface, however, is a specular reflector, and reflection of radiation from it follows the geometrical law of reflection. This geometric law requires that the angle with respect to the normal to the surface of the reflected ray equal the angle of incidence of the incident ray and that the reflected ray be in the same plane as the incident ray. Specular reflection is depicted in Figure 1. Reflectance, ρ_s , of the water surface is given by the Fresnel equation

$$\rho_s = \frac{1}{2} \left[\frac{\sin^2(\theta_i - \theta_r)}{\sin^2(\theta_i + \theta_r)} + \frac{\tan^2(\theta_i - \theta_r)}{\tan^2(\theta_i + \theta_r)} \right], \quad (1)$$

where θ_i is the incident angle and θ_r is the angle of refraction.

The transmitted part of the incident ray experiences refraction at the water surface, as shown in Figure 2. Snell's law, given by

$$\frac{\sin(\theta_i)}{\sin(\theta_r)} = n, \quad (2)$$

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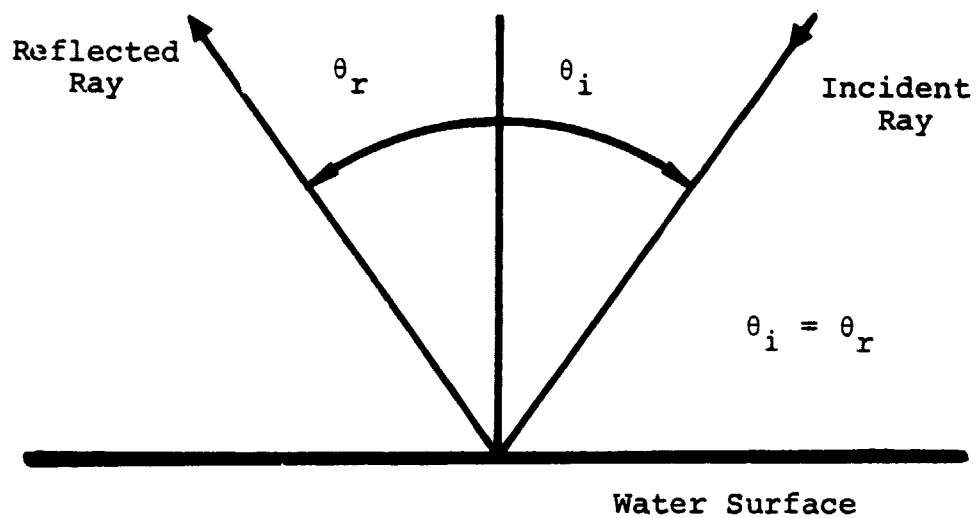


FIGURE 1. SPECULAR REFLECTION FROM A SMOOTH WATER SURFACE.
THE INCIDENT AND REFLECTED RAYS ARE IN THE
SAME PLANE.

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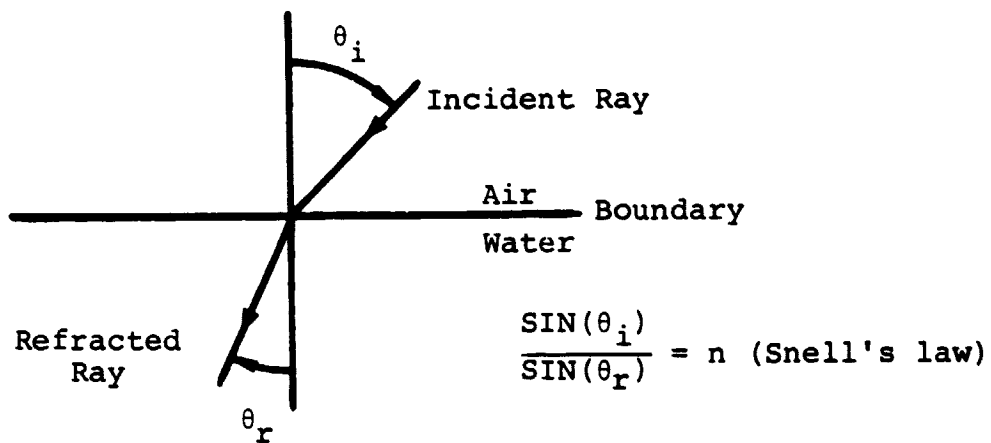


FIGURE 2. REFRACTION OF THE TRANSMITTED BEAM AT THE
WATER SURFACE AND THE LAW OF REFRACTION
 n = REFRACTIVE INDEX OF WATER RELATIVE TO AIR.

describes the relationship between the angle of incidence of the incoming beam, θ_i , the angle with respect to the normal to the surface of the refracted beam, θ_r , and the refractive index of air relative to water. The refracted beam lies in the same plane as the incident beam. The phenomenon of refraction is depicted in Figure 2. The law of refraction requires that for a water surface, downward sky radiation and direct sunlight enter the water within 48.5° of the vertical. Only when the water surface is roughened by wind or another disturbance can direct sunlight or sky radiation penetrate the water surface outside this range of angles. Back-scattered radiation from beneath the water-air interface also experiences refraction on reaching the water surface. When the water surface is calm, upward radiation incident at angles greater than 48.5° with the vertical is totally internally reflected [2]. Thus, downward radiation beneath the water surface at angles with the vertical greater than 48.5° is upward radiation in the water which has been totally internally reflected [3].

Equations 1 and 2 show that the reflectance and transmittance of the water surface are dependent on the refractive index of water. The refractive index is influenced by changes in temperature and by the concentration of various solutes in the water. Figure 3 shows how the reflectance function, equation 1, varies as a function of the angle of incidence for refractive indices of 1.20, 1.33, 1.40, 1.45. This range of refractive indices encompasses the range of natural variability in the refractive index for water; and Figure 3 shows that, over this range, variation in the refractive index is of little importance in determining the surface reflectance. For all of the calculations shown in this report, a refractive index of $4/3$ is used. Figure 4 shows transmittance, T , and reflectance, ρ_s , as a function of angle of incidence of the incoming radiation for a refractive index of 1.33.

Since many applications of remote sensing over water require

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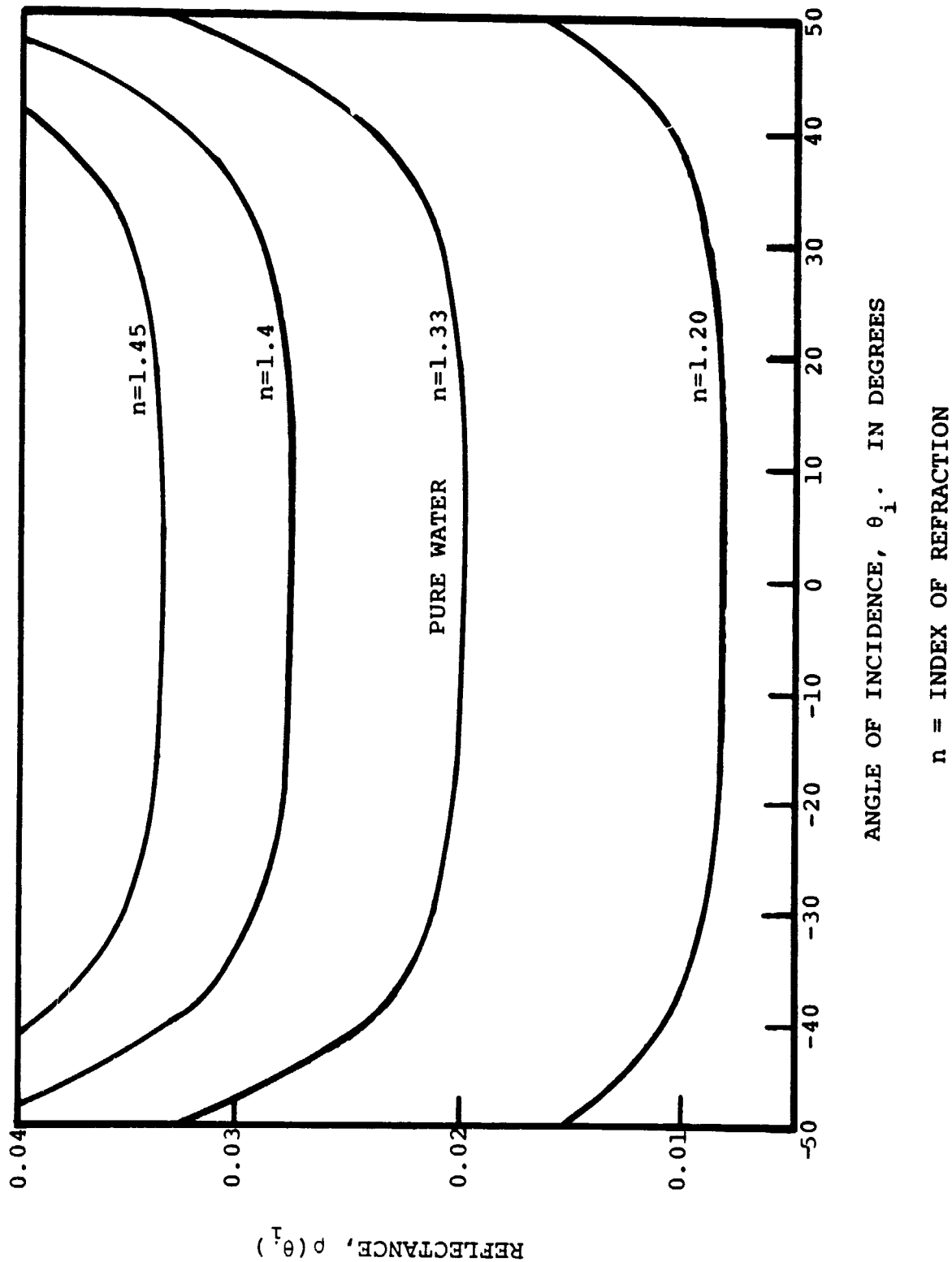


FIGURE 3. FRESNEL REFLECTANCE FOR UNPOLARIZED LIGHT AS A
FUNCTION OF ANGLE OF INCIDENCE OF INCOMING BEAM.

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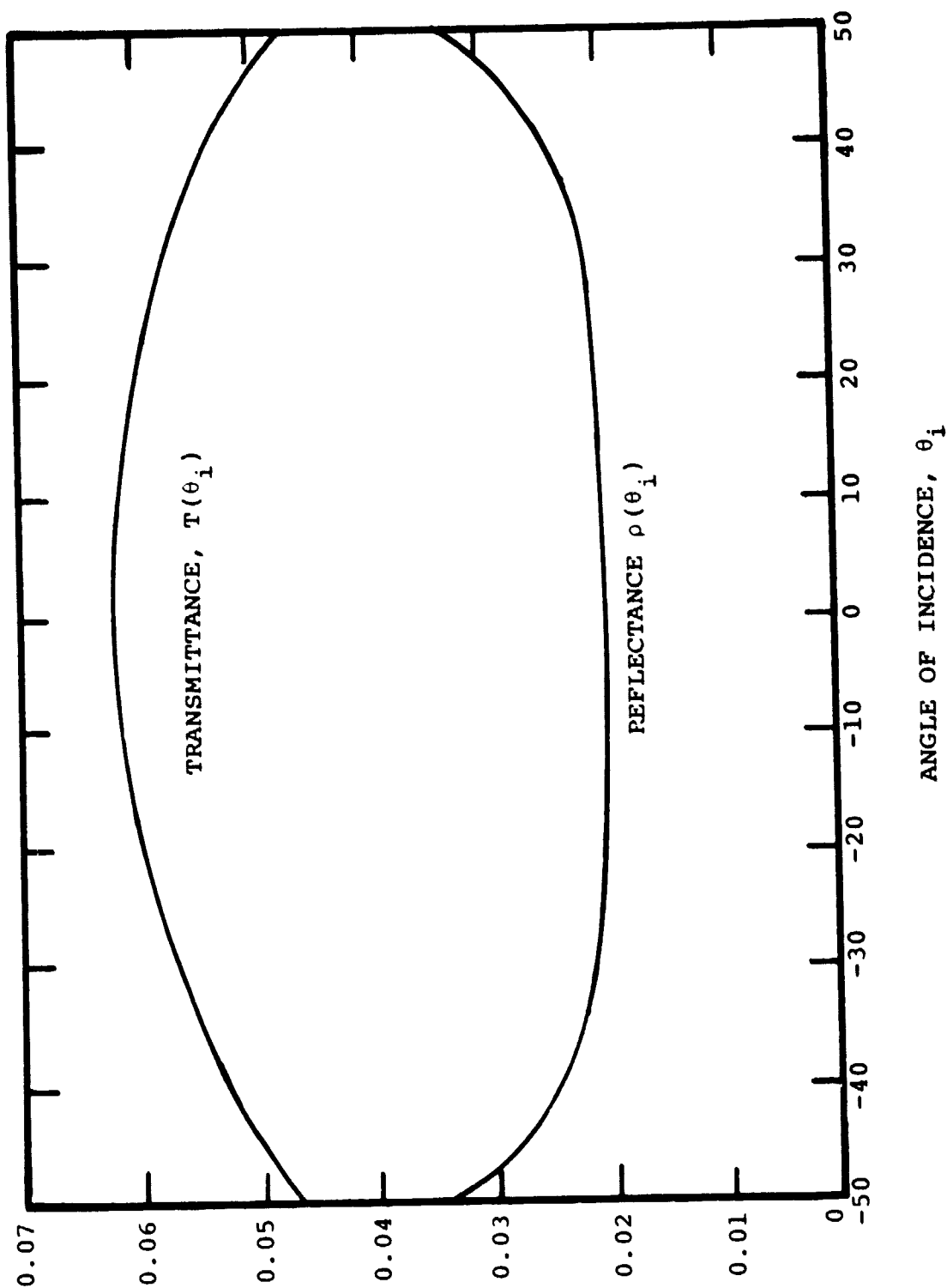


FIGURE 4. FRESNEL REFLECTANCE AND TRANSMITTANCE AS A FUNCTION OF ANGLE OF INCIDENCE OF INCOMING BEAM. REFRACTIVE INDEX = 1.33.

examining the signature from beneath the water surface, Fresnel reflectance is of paramount importance. Figure 4 shows that transmittance is highest when the incident beam is perpendicular to the water surface, while the reflectance is lowest at this angle. Thus, from a consideration of the Fresnel formulas, one would expect the return from beneath the water surface relative to the total, to be greatest when viewing that surface at the nadir. Contributing to this effort is the fact that most scattering phase functions for polydispersions in water have a secondary peak at 180° [2]. Figure 4 shows that while the Fresnel reflectance function is at a minimum at the nadir view angle, this function changes very little out to viewing angles as great as 40° , at which point it begins to rise steeply. Internal reflection at the water surface of upwelling radiation is also at a minimum normal to the water surface. Thus, for remote sensing work where the return from beneath the water surface is of greatest interest, scan angles should be maintained within 40° to 45° of the vertical. Beyond an angle of 48.5° no radiation from beneath the water surface will reach the sensor when the water surface is calm. Scanning in the solar plane is also problematic in this regard since specularly reflected light on the solar side of the scan plane would saturate the sensor. For some purposes, such as viewing the glitter pattern on the water surface, scanning in the solar plane may be desirable.

3.2 ABSORPTION AND SCATTERING

Water is a good absorber of electromagnetic radiation. Only in the relatively narrow spectral region from about 400 to 600 nanometers is the transparency of water such that radiation can penetrate more than a few meters in depth below the water surface. Both at wavelengths shorter than 400 nm and longer than 600 nm, absorption increases rapidly and only very small amounts of radiation are scattered back out of the water

into the atmosphere. At the very short wavelengths this radiation is further strongly attenuated in the atmosphere.

Scattering of radiation in water is caused by water molecules, by dissolved salts and by particles in suspension. These effects are usually assumed to be additive [4].

Scattering by water molecules is described by fluctuation theory which predicts scattering of radiation as a result of molecular movements which cause fluctuations in the density of the medium. As in Rayleigh scattering, this type of scattering is proportional to λ^{-4} , where λ is the wavelength of the radiation being scattered. The effect of dissolved salts on the molecular scattering phase function is usually small enough to be neglected.

In general, most scattering in water is accomplished by particles in suspension [4]. Particulate matter in water derives from runoff from land, deposition from the atmosphere, and organic processes within the water. Thus, particles may be quite irregular in shape and particle size distributions are difficult to characterize precisely [2]. Because some sources of particles such as runoff of organic processes may be highly localized in space, size distributions may vary greatly in space and time. Although particle shapes vary considerably from the spherical ideal of scattering theory, it has been shown [5,6] that systems of irregularly shaped particles can be adequately approximated by systems of polydisperse systems of spherical particles. The major observed features of phase functions of particulate suspensions in water are a strong forward scattering peak, a broad minimum around 100° - 130° and a small secondary peak in the back scattering direction [4].

3.3 OPTICAL PROPERTIES OF A WIND ROUGHENED WATER SURFACE

Roughness of the water surface caused by wind presents an additional problem in the calculation of the optical properties of the surface. Waves increase the angle of incidence of direct

radiation for high solar elevations. The effect on the Fresnel reflectance, however, is of little consequence since the reflectance does not vary much with solar zenith angle for zenith angles less than 40° (see Figure 3). Waves reduce the angle of incidence of direct radiation from a low sun, greatly reducing the reflectance of the water surface. Cox and Munk [7] have shown that wave action becomes a significant factor for solar elevations below 20° . At these low sun elevations, reduced reflection, shadowing and multiple reflections greatly reduce the reflected radiance.

The reflection of diffuse radiation by the water surface is little affected by surface roughness [2], although complete agreement on this matter is lacking [8]. Burt [9] found that the albedo of a wind roughened water surface was slightly less than the albedo of a smooth water surface--a decrease from 6.6% to 5.7% for the roughened surface. Cox and Munk [7] measured a small increase in the albedo of a smooth water surface of 5% to 5.5% for a water surface roughened by waves. Kondratyev [8] on the other hand, calculates that where the solar zenith angle is 0° the albedo of calm water surface of 2.1% will increase to 13.1%. When the solar zenith angle is 30° , the increase will be from 2.2% to 3.8%, and for a solar zenith angle of 60° there will be a decrease from 6.2% down to 2.4% for a roughened surface. Plass et al. [10], using a Monte Carlo model of the atmosphere ocean system, demonstrate that the downward flux just below the surface always increases with wind speed, even at high sun elevations. They attribute this result to the fact that more sky radiance near the horizon enters the water when waves are present.

The effect of waves on the radiance of the water surface can be calculated if the probability distribution of surface slopes is known. For an observer looking down on a water surface, the specular angle will vary from place to place over the surface of the water. Since in most remote sensing applications

the light source (Sun) and observer (sensor) are high enough above the surface and the region viewed sufficiently small that variation in the specular angle can be neglected. The radiance of the surface is then directly proportional to the probability of finding a surface element with slope, S_0 , at the specular angle [11]. If p is this probability, the radiance of the surface, L_s , at vertical angle i is given by

$$L_s(i) = L_0(i) \rho_f(i) p \quad (3)$$

where

$L_0(i)$ is radiance incident at the surface at vertical angle i , and

$\rho_f(i)$ is the Fresnel reflectance at vertical angle i .

Duntley [12] and Cox and Munk [7,13] have studied the statistical distribution of wave slope as a function of wind speed. Observations of the effect of wind speed on spatially or temporally averaged reflectance of the water surface indicate that it is not significant for view angles less than 70° from vertical. Angles in excess of 50° from the vertical are seldom used in remote sensing systems because of the large optical air mass at these angles.

COMPONENTS OF REFLECTED AND PATH RADIANCE IN REMOTE SENSING OVER WATER

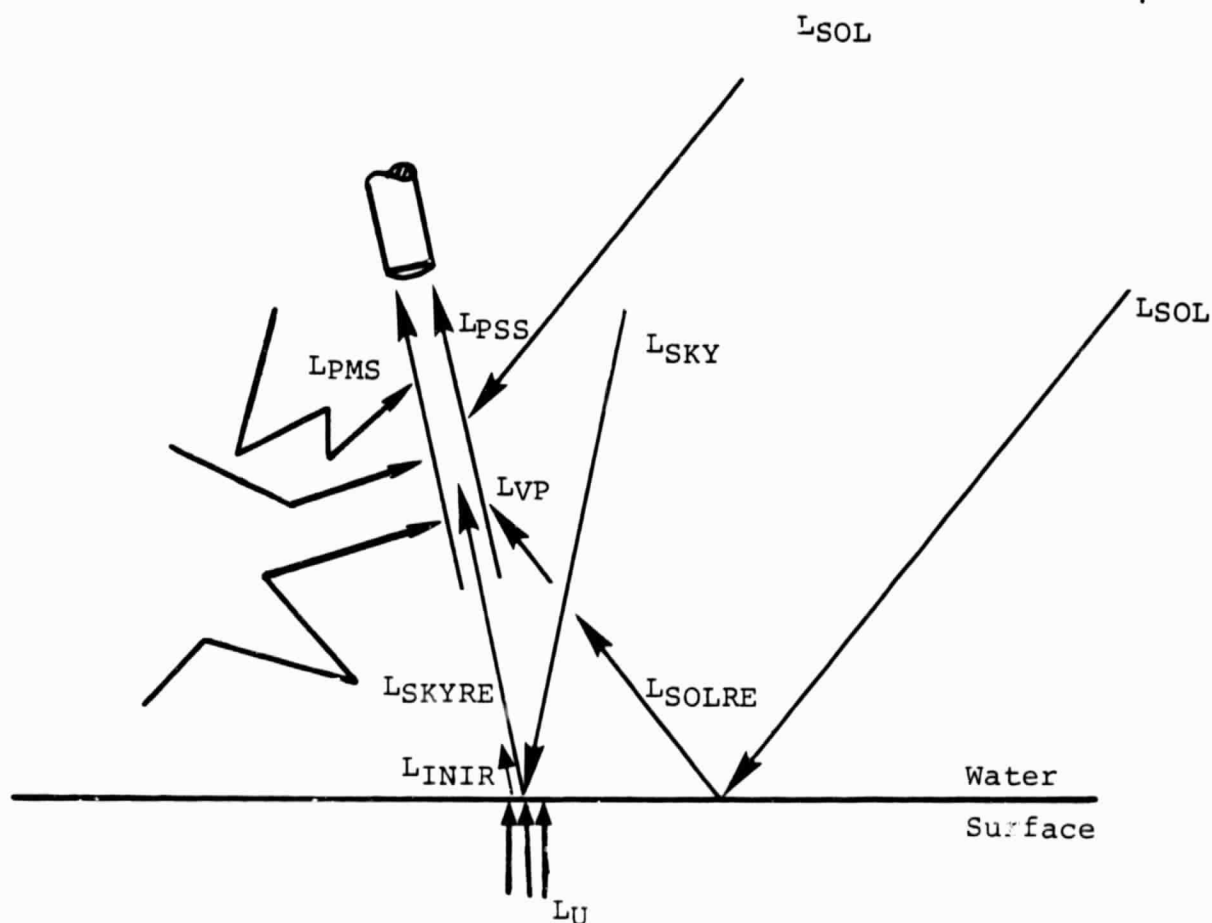
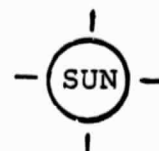
In many applications of oceanographic remote sensing the quantity of greatest interest is the radiance information transmitted from below the water surface to a sensor, sometimes called the intrinsic radiance. To determine this quantity from raw remote sensing data we must not only estimate atmospheric path radiance but also the magnitude of radiance reflected off the water surface and transmitted to the sensor. In this chapter we describe in detail analytical models appropriate for estimating the following quantities: reflected sky radiance, $LSKYRE$; singly scattered reflected solar path radiance, Lvp ; singly scattered path radiance, $Lpss$; and multiply scattered path radiance, $Lpms$. The first two of these quantities are radiances resulting from specular reflection off the water surface, the latter two are atmospheric path radiances. Each of these radiances augments the radiance detected by a sensor, masking the radiance signal from beneath the water surface, as shown in Figure 5.

4.1 REFLECTED SKY RADIANCE, $LSKYRE$

The geometry for sky radiance reflected into the line of sight of the sensor is depicted in Figure 6, where θ represents the nadir view angle of the sensor, ϕ the azimuth of the sensor scanning plane. We assume a plane parallel uniform atmosphere. Sky radiance downwelling in the scan plane and incident at an angle θ with the normal to the surface is reflected in the direction of the sensor, and attenuated by the atmosphere as it travels to the sensor. The surface reflectance is given by the Fresnel formula for unpolarized light described in Chapter 3.

We consider in this model only singly scattered sky radiance, generated by scattering of solar beam radiation along the straight

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- L_{SOL} = direct beam radiance;
- L_{PMS} = multiply scattered path radiance;
- L_{PSS} = singly scattered path radiance;
- L_{SKY} = sky radiance;
- L_{SKYRE} = reflected sky radiance;
- L_{SOLRE} = reflected solar radiance;
- L_U = upward scattered radiance beneath water surface;
- L_{INIR} = radiance scattered from beneath water surface.

FIGURE 5. COMPONENTS OF TOTAL RADIANCE DETECTED BY SENSOR

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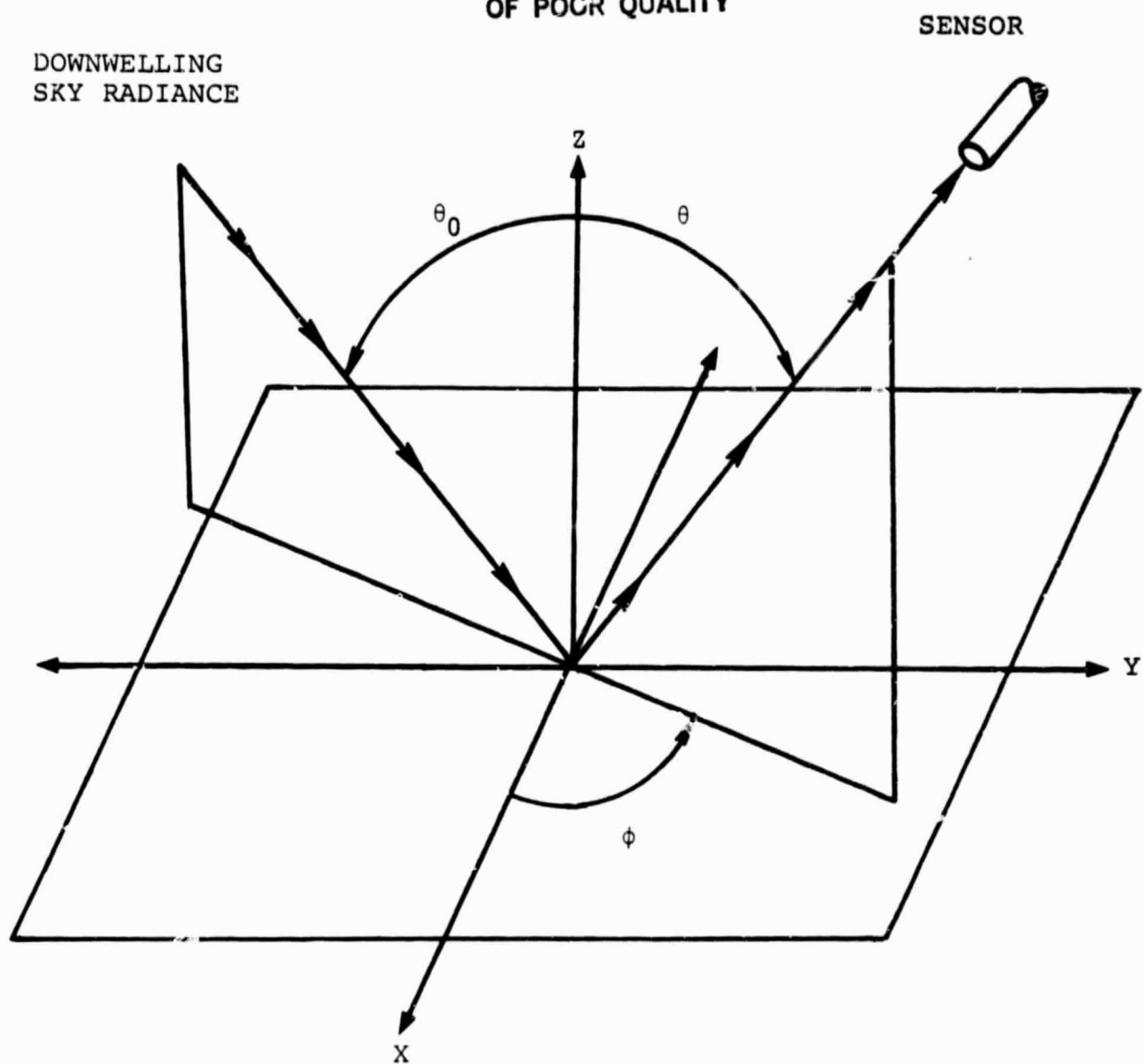


FIGURE 6. GEOMETRY FOR REFLECTED SKY
RADIANCE COMPUTATION.

line path from the top of the atmosphere to the water surface. In Figure 6, this straight line path has direction $(-\mu, \phi)$, where $\mu = \cos\theta$. In order to define the angle of scattering we must first define two vectors, one defining the direction of a photon leaving the Sun and the other the direction of the singly scattered sky radiation. If $(-\mu_0, \phi_0)$ is the direction of the photon leaving the Sun (where μ_0 is the cosine of the solar zenith angle, θ_0 , and ϕ_0 is the photon azimuth), the vector direction of the photon leaving the Sun is

$$\hat{L}_{\text{SOL}} = \begin{bmatrix} \sin(\pi-\theta_0)\cos\phi_0 \\ \sin(\pi-\theta_0)\sin\phi_0 \\ \cos(\pi-\theta_0) \end{bmatrix} = \begin{bmatrix} \sqrt{1-\mu_0^2} \cos\phi_0 \\ \sqrt{1-\mu_0^2} \sin\phi_0 \\ -\mu_0 \end{bmatrix}, \quad (4)$$

and

$$\hat{L}_{\text{SKY}} = \begin{bmatrix} \sin(\pi-\theta)\cos\phi \\ \sin(\pi-\theta)\sin\phi \\ \cos(\pi-\theta) \end{bmatrix} = \begin{bmatrix} \sqrt{1-\mu^2} \cos\phi \\ \sqrt{1-\mu^2} \sin\phi \\ -\mu \end{bmatrix}. \quad (5)$$

The cosine of the scattering angle, χ_{SKY} , is given by the dot product $\hat{L}_{\text{SOL}} \cdot \hat{L}_{\text{SKY}}$, i.e.

$$\cos\chi_{\text{SKY}} = \mu \mu_0 + \sqrt{1-\mu^2} \sqrt{1-\mu_0^2} \cos(\phi-\phi_0). \quad (6)$$

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The equation for singly-scattered sky radiance, L_{SKYSS} , is the well-known formula [14]

$$L_{\text{SKYSS}} = \frac{\omega_0 \mu_0 E_0 P(\chi_{\text{SKY}})}{4\pi(\mu_0 - \mu)} \left[e^{-\tau_0/\mu_0} - e^{-\tau_0/\mu} \right] \quad (7)$$

where E_0 = solar irradiance at the top of the atmosphere;
 ω_0 = atmospheric single scattering albedo;
 τ_0 = optical thickness of atmosphere;
 $P(\chi_{\text{SKY}})$ = scattering phase function for scattering angle χ_{SKY} .

When the sky radiance is reflected off the surface of the water it is diminished by the Fresnel reflectance of the water surface, ρ_F , and further attenuated by the atmosphere on its way to the sensor. Thus, the complete formula for the reflected sky radiance is

$$L_{\text{SKYRE}} = \rho_F e^{-(\tau_0 - \tau)/\mu} L_{\text{SKYSS}} \quad (8)$$

where μ is the cosine of the scan angle, τ is the optical depth of the sensor, and $e^{-(\tau_0 - \tau)/\mu}$ is the transmittance of the atmosphere between the water surface and the sensor. At this point we take note of the fact that in the above formula for the singly scattered reflected sky radiance, the reflectance of the water surface and the transmittance of the atmosphere are opposing effects. Assuming the refractive index of water to be 1.33, the

Fresnel reflectance of water at the nadir view angle reaches a minimum of 0.021 and attains a maximum of 1.0 at the grazing angle. Typical values of ρ_F for angles commonly used in remote sensing range between 2.1 at the nadir to 3.0 at a scan angle of 46° . The transmittance, on the other hand, reaches a maximum at the nadir and becomes increasingly small as the scan angle increases. These effects will be discussed further in the following chapter.

4.2 SINGLY SCATTERED REFLECTED SOLAR RADIANCE, L_{VP} .

The phenomenon of specular reflection produces an image of the radiation source on the surface of the water. We refer to the image of the Sun on the water surface as the virtual Sun. If the scan plane is coincident with the solar plane and the sensor is scanning on the solar side of the scan plane at a view angle equal to the solar zenith angle, the field of view becomes saturated with the radiance of the Sun's image. Radiance from the virtual Sun is also scattered into the line of sight of the sensor. In the terminology of this report, we refer to singly-scattered path radiance from the Sun as virtual Sun path radiance, L_{VP} .

To find the scattering angle for the computation of singly-scattered virtual Sun path radiance, we note that the zenith angle of a photon leaving the virtual Sun is θ_{SUN} and the azimuth angle is $\phi_0 = \phi_{SUN} + \pi$ (see Figure 7). The vector direction of a photon leaving the virtual Sun is

$$\hat{L}_{VSOL} = \begin{bmatrix} \sin\theta_{SUN} \cos\phi_0 \\ \sin\theta_{SUN} \sin\phi_0 \\ \cos\theta_{SUN} \end{bmatrix} = \begin{bmatrix} \sqrt{1-\mu_0^2} \cos\phi_0 \\ \sqrt{1-\mu_0^2} \sin\phi_0 \\ \mu_0 \end{bmatrix} \quad (9)$$

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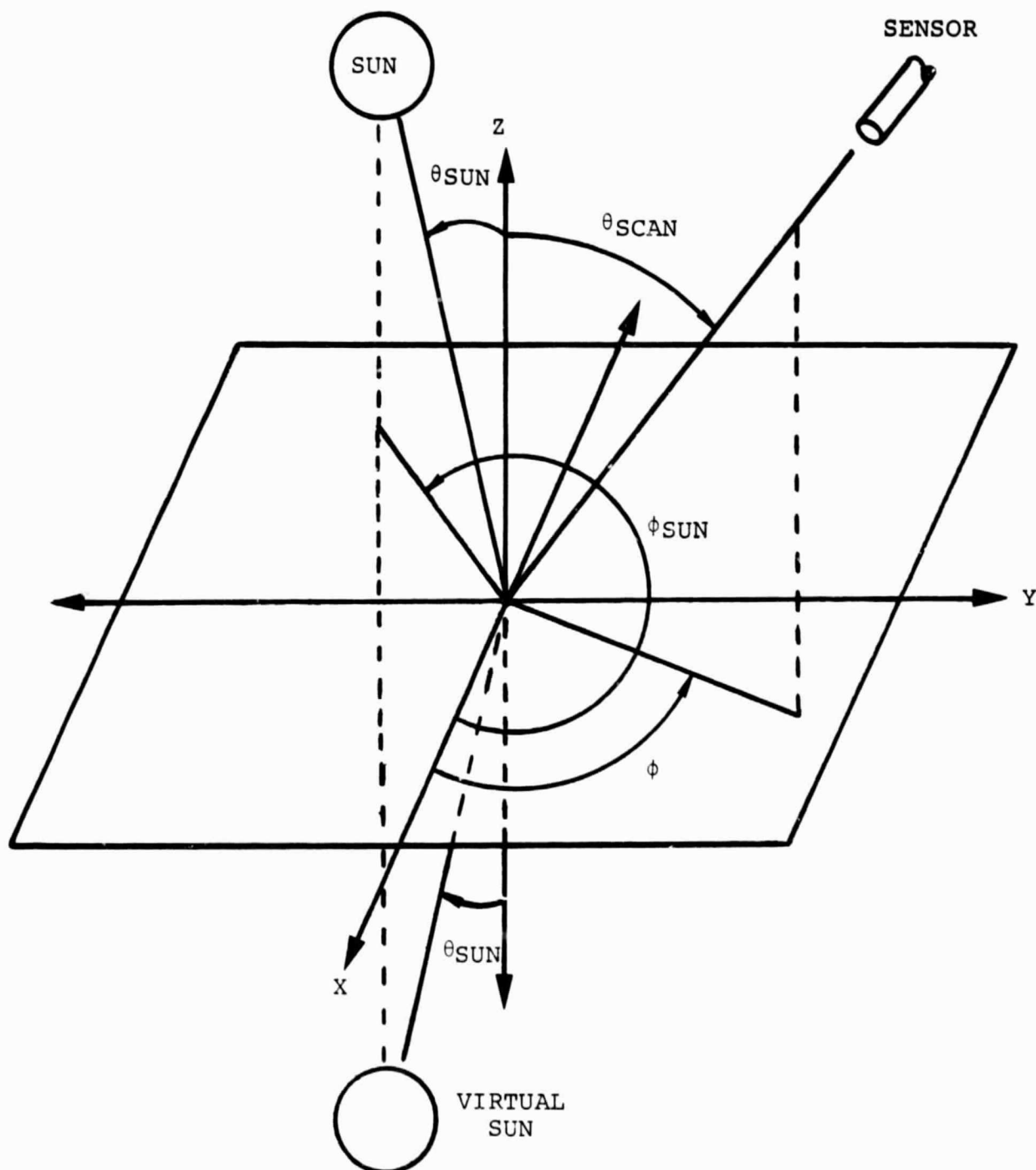


FIGURE 7. GEOMETRY FOR VIRTUAL SUN PATH RADIANCE CALCULATION.

and the vector direction into which photons from the virtual Sun are scattered, creating virtual Sun path radiance is

$$L_{VP} = \begin{bmatrix} \sin\theta\cos\phi \\ \sin\theta\sin\phi \\ \cos\theta \end{bmatrix} = \begin{bmatrix} \sqrt{1-\mu^2} \cos\phi \\ \sqrt{1-\mu^2} \sin\phi \\ \mu \end{bmatrix} . \quad (10)$$

The scattering angle is the dot product,

$$L_{VSOL} \cdot L_{VP} = \mu \mu_0 + \sqrt{1-\mu^2} \sqrt{1-\mu_0^2} \cos(\phi-\phi_0) . \quad (11)$$

As in the case of singly scattered sky radiance, the same types of physical interactions which generate singly-scattered sky radiance from direct solar radiation also scatter radiation from the virtual Sun to generate virtual Sun path radiance. Thus, we may use the same equation for singly-scattered sky radiance, with some modifications, to find the singly-scattered virtual Sun path radiance, L_{VP} . One difference is in that the computation of L_{VP} we will now sum the scattered radiation over a path beginning at $\tau = \tau_0$ (the optical depth of the scene viewed by the sensor) and ending at the optical depth of the sensor, τ . If we denote the irradiance of the virtual Sun by E_0' , we obtain the following formula for singly-scattered path radiance from the virtual Sun:

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$$L_{VP} = \frac{\omega_0 \mu_0 E_0' p(\cos \chi_{VP})}{4\pi(\mu_0 - \mu)} \left[e^{-(\tau_0 - \tau)/\mu_0} - e^{-(\tau_0 - \tau)/\mu} \right] . \quad (12)$$

We define E_0' by noting that the E_0' is the image of the Sun reflected in the water surface. Hence, the irradiance of the virtual Sun is the irradiance of the true Sun at the top of the atmosphere attenuated by the atmosphere and the reflectance of the water surface, i.e.,

$$E_0' = \rho_F e^{-\tau_0/\mu_0} E_0 \quad (13)$$

where ρ_F is the Fresnel reflectance and $e^{-\tau_0/\mu_0}$ is the transmittance of the atmosphere.

4.3 SINGLY SCATTERED PATH RADIANCE, L_{PSS} .

The geometry for singly-scattered path radiance is shown in Figure 8. The formula for singly-scattered path radiance is similar to that for singly-scattered path radiance from the Sun--the same straight line path from τ_0 to τ is used, but E_0 is substituted for E_0' in the formula. The cosine of the scattering angle for singly-scattered path radiance, $\cos \chi_{PSS}$ is also the negative of the cosine of the scattering angle used to compute the phase function for L_{VP} . Thus, $\cos \chi$ is the dot product, $\hat{L}_{SOL} \cdot \hat{L}_{VP}$, vector directions which have already been defined. The formula used to compute singly-scattered solar path radiance is

$$L_{PSS} = \frac{\omega_0 \mu_0 E_0 p(\cos \chi_{PSS})}{4\pi(\mu + \mu_0)} e^{-\tau_0/\mu_0} \left[e^{(\tau_0 - \tau)/\mu_0} - e^{-(\tau_0 - \tau)/\mu} \right] \quad (14)$$

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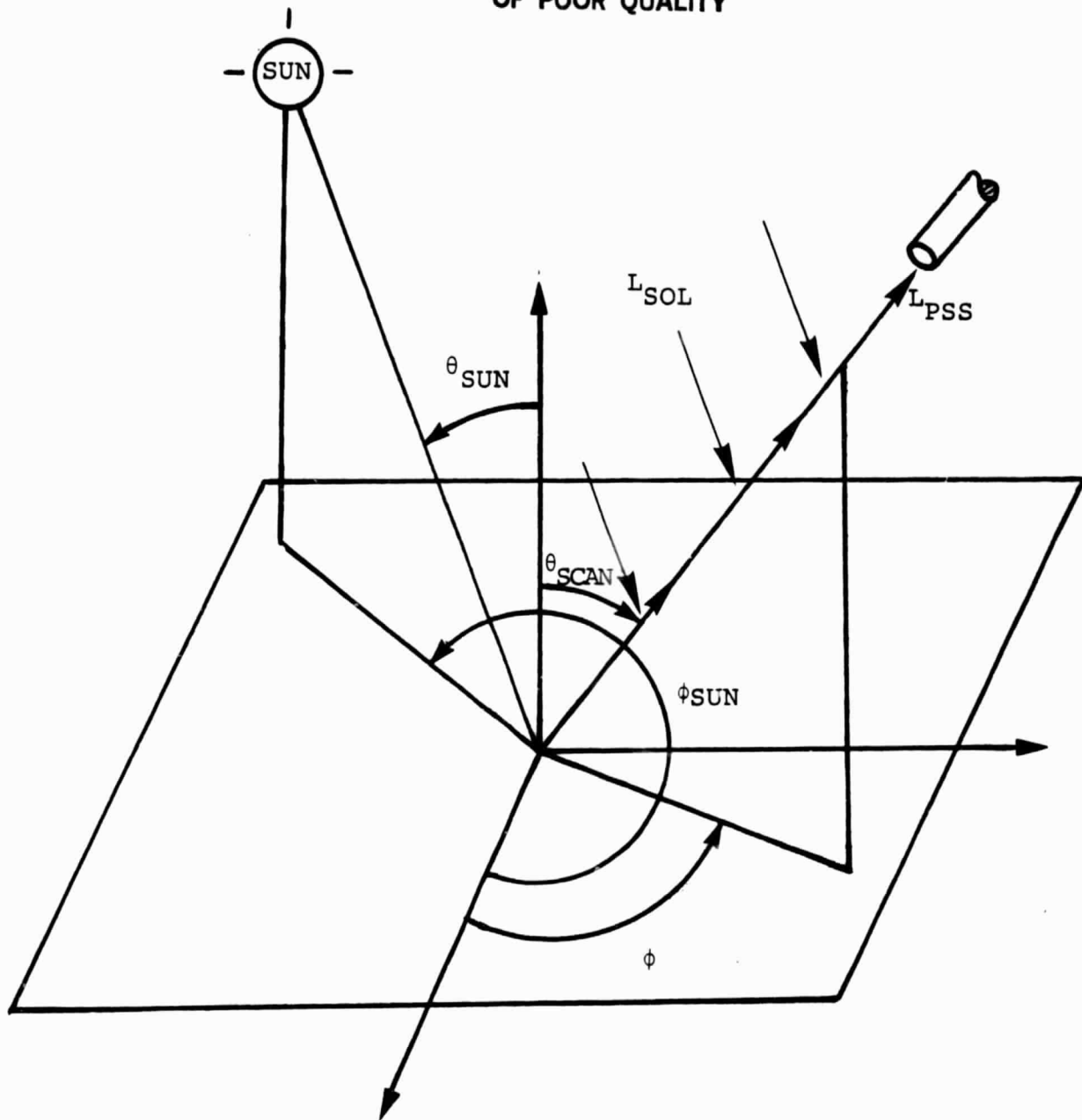


FIGURE 8 . GEOMETRY FOR SINGLY SCATTERED PATH RADIANCE, L_{PSS} , CALCULATION.

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where the variables in the above equation are as previously defined.

4.4 MULTIPLY SCATTERED PATH RADIANCE, L_{PMS} .

To compute multiply scattered path radiance we use an analytical approximation described in detail in an earlier report [15]. The formula for the computation of L_{PMS} is

$$\begin{aligned}
 L_{PMS} = & \frac{E_0}{4\pi [\mu_0 + (1-\eta)\tau_0]} \\
 & \left(\left\{ (1-\eta)\tau_0 [p(\mu, \phi, \mu_0, \pi + \phi_0) + p(\mu, \phi, -\mu_0, \phi_0)] + \mu_0 p(\mu, \phi, -\mu_0, \phi_0) \right. \right. \\
 & + \frac{2\mu_0^2 \rho}{1+2(1-\eta)(1-\rho)\tau_0} \left. \left\{ \left[1 - e^{-(\tau_0 - \tau)/\mu} \right] + \left\{ (1-\eta) p(\mu, \phi, \mu_0, \pi + \phi_0) \right. \right. \right. \\
 & + p(\mu, \phi, -\mu_0, \phi_0) \left. \left. \left. - \frac{8(1-\eta)\mu_0^2 \rho}{1+2(1-\eta)(1-\rho)\tau_0} \right\} \left[(\tau_0 + \mu) e^{-(\tau_0 - \tau)/\mu} - (\tau + \mu) \right] \right\} \right)
 \end{aligned}
 \tag{15}$$

The single-scattering phase functions are given by:

$$p(\mu, \phi, \mu_0, \pi + \phi_0) = p \left[\mu \mu_0 - \sqrt{(1-\mu^2)(1-\mu_0^2)} \cos(\phi - \phi_0) \right]$$

$$p(\mu, \phi, -\mu_0, \phi_0) = p \left[-\mu \mu_0 + \sqrt{(1-\mu^2)(1-\mu_0^2)} \cos(\phi - \phi_0) \right].$$

COMPUTER MODEL AND RESULTS

In this chapter we describe implementation of the formulas discussed in Chapter 4 in the computer program called ATCOR. Many of the details of ATCOR have been presented in a previous report [16], so only a brief description of the entire program will be given here. Our discussion will focus primarily on the subroutine ATMSFR, in which the formulas presented in Chapter 4 have been introduced.

5.1 SPECIFICATION OF SOLAR AND SENSOR GEOMETRY

The geometric relationship of the sensor to the environment is shown in Figure 9. The geographic coordinates of the sensor locate the center of a spherical coordinate system used to define angles needed in model calculations. In the diagram the scanner scans along a path from P_1 (the first pixel) to P_n (the last pixel). The azimuth of the scan plane is measured in the counterclockwise direction from north to the first pixel and is read into the program by the routine ATMSFR. The first pixel is always 90° in a clockwise direction from a vector pointing in the direction of the flight.

The solar zenith and azimuth angles are computed automatically once the latitude, longitude, date, time of day (standard time), and zone number are specified. The extraterrestrial solar irradiance is also computed based on these inputs.

5.2 ATMOSPHERIC CORRECTION OPTIONS

Two input parameters set by the user determine which calculations are performed in routine ATMSFR. These parameters are SCATT and OPTION. SCATT may assume the value of either 0 or 1; OPTION can take on the values of 1, 2, or 3. If SCATT is 0, only multiply scattered path radiance is calculated and the value of OPTION can be any integer and will be ignored since only LPMS is then calculated.

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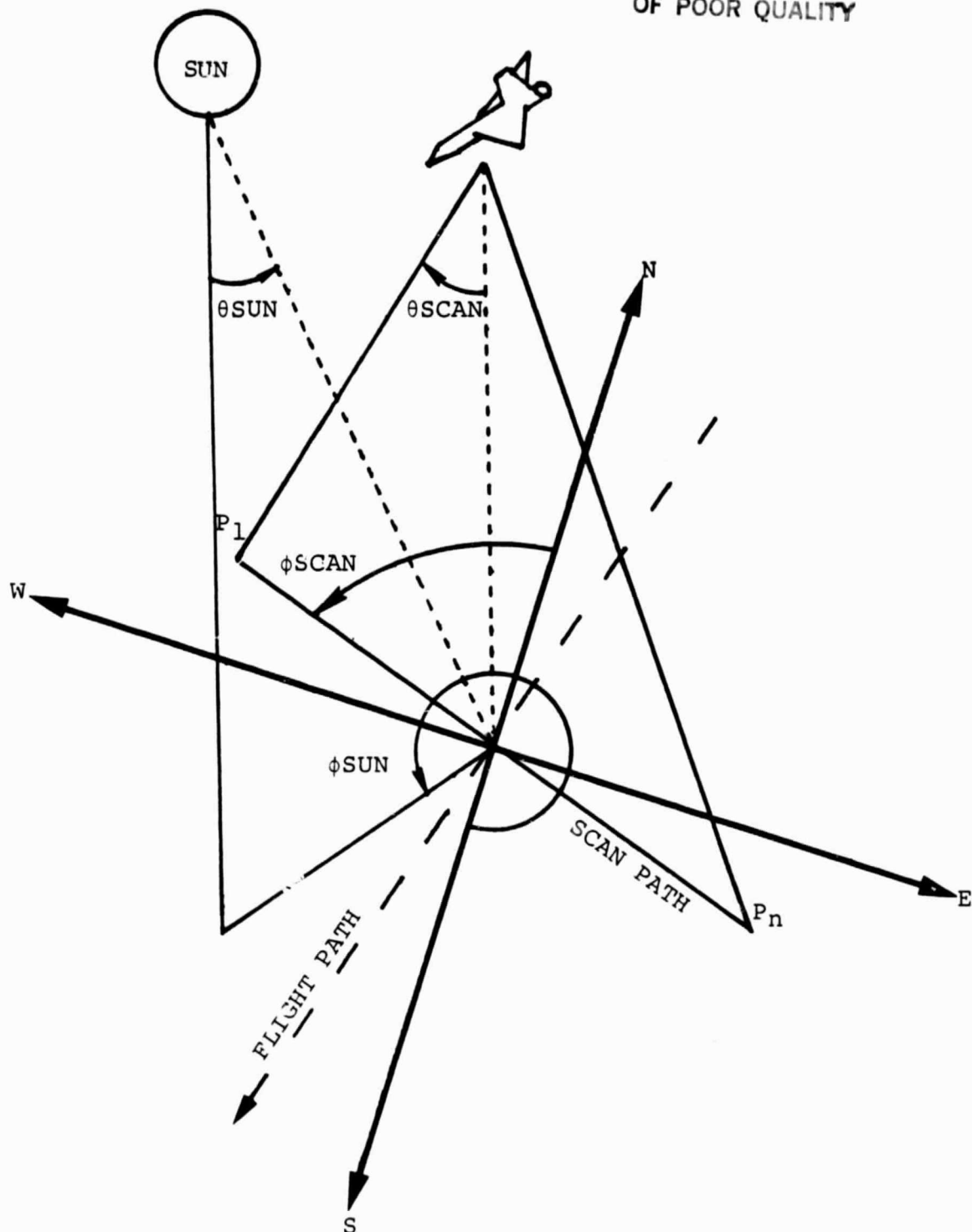


FIGURE 9. SCAN PLANE AND SOLAR GEOMETRY. ϕ_{SCAN} IS MEASURED COUNTERCLOCKWISE FROM NORTH TO PIXEL #1, ϕ_{SUN} IS MEASURED COUNTERCLOCKWISE FROM SOUTH. POSITIVE SCAN ANGLES ARE MEASURED FROM THE FIRST PIXEL TO THE NADIR. NEGATIVE SCAN ANGLES ARE MEASURED FROM THE NADIR TO THE LAST PIXEL. P_1 AND P_n ARE THE FIRST AND LAST PIXELS, RESPECTIVELY.

If SCATT is 1, single scattering computations are performed and the value of OPTION is used to determine which values to calculate. If OPTION = 1, singly scattered path radiance, LPSS, and reflected sky radiance, LSKYRE, are computed. If OPTION = 2, LPSS and virtual sun path radiance, LVP, are computed. If OPTION = 3, LPSS, LSKYRE and LVP are calculated.

OPTION and SCATT are read into the data file on logical unit 4. The format for this record is (5I5), and the variables read in are:

FSTP LSTP PTINC SCATT OPTION

where

FSTP = the number of the first pixel to be processed,
LSTP = the number of the last pixel to be processed,
PTINC = the pixel increment to use in the processing,
and SCATT and OPTION are as previously defined.

The output file which is used by ATMSFR is given in Table 1 and the new subroutine ATMSFR2 is given in Table 2.

5.3 MODEL INPUT PARAMETERS

In addition to the geometric parameters, we must specify parameters characterizing the medium and the measurement system.

The model makes use of several "altitude" values which must be input by the user. First, one must know the actual altitude (km) of the sensor above the surface. Second, one must know the pressure (millibars) of the atmosphere at the surface, and third, one must know the atmospheric pressure (in millibars) at flight altitude. If only the altitudes are known, one can use the tables relating pressure to altitude as given by the U.S. Standard Atmosphere [17].

TABLE 1
INPUT FILE ON LOGICAL UNIT 4
FOR USE BY ATMSFR

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LINE NO	READ OCCURS IN ROUTINE	INPUT VARIABLES	FORMAT
1	ATMSFR	FSTP,LSTP,PTINC, SCATT,OPTION	(5I5)
2	DATE	MNTH,DAY,YEAR	(3I4)
3	ANGLES	HOURL,MIN,SEC	(2I5,F6.3)
4	ANGLES	NZ	(I5)
5	ANGLES	LATD,LATM,LATS,LONGD, LONGM, LONGS	(2I5,F6.3,2I5, F6.3)
6	ATMSFR	ZSCAN,ZGRND,LSW	(2F8.5,I5)
7	ATMSFR	(WAVE(I),I=1,QNCHAN)	(10F8.5)
8	ATMSFR	(RHO(I),I=1.QNCHAN)	(10F8.5)
9	PHASE	R,IM	(2F8.6)
10 11-14	PHASE	NWT,NANG,(C(I),I=1,NANG)	(2I5/(10F8.6))
15-79	PHASE	(WTAB(I),(PF(I,J), J=1,NANG),I=1,NWT)	(F10.6/(10 F8.4))
80	RAYLEI	PRESO,PRESZ	(2F10.4)
81	OZONE	NOZ,NPROF,NO3W1,NO3W2	(4I5)
82	OZONE	WAVC1,WAVC2	(2F8.4)
83-153	OZONE	((ZOZ(IZ),O3INT(IZ,IP), IZ=1,NOZ),IP=1,NPROF)	(F7.0, E11.4)
154	OZONE	(O3MAX(IP),IP=1,NPROF)	(10E13.6)
155-169	OZONE	(WAVO3(I),A(I),I=1,NO3W) (NOTE:NO3W=NO3W1 + NO3W2)	(F7.0,E11.4)

TABLE 1 (Cont.)
INPUT FILE ON LOGICAL UNIT 4
FOR USE BY ATMSFR
(CONTINUED)

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LINE NO.	READ OCCURS IN ROUTINE	INPUT VARIABLES	FORMAT
170	OZONE	NOP	(4I5)
171	THICK	NTEX	(I2)
177	THICK	(WAVEX(I),TAUEX(I),I=1,NTEX)	(2F8.4)
178	PARAMS	FSCAT	(10F8.6)
179	AERO	NAER,MPROF,NUZ,MAXG	(5I5)
180	AERO	I PROF	(I5)
181-184	AERO	(WAVAER(I),RIN(I),I=1,NAER)	(10F8.4)
185-230	AERO	((ZUN(IZ),UNIZ(IZ,IP),IZ=1, NUZ),IP=1,MPROF)	(F7.0,E11.4)
231-258	AERO	(X(I),Z(I),I=1,MAXG)	(F7.0,E11.4)
259	ATMSFR	PHID,PHIM,PHIS	(2I3,F6.3)
		\$ENDFILE	

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TABLE 2 (Cont.)

61	C	COMMON /QCOM/ QCVERS(4), QNSAS(6), QNSS, QLINE, QSLIME,
62	C	QSTEP, QNCHAN, QNSS, QNCHAN, QNDACH, QFILE,
63	C	QBNONE, QBPAST, QBPOLY, QBTASE, QFLDID(24), QSPARE(93),
64	C	QBLSEG, QBGEOB, QBLSEQ, QBIANG, QDANG, QREAL, QFACT,
65		QTTITLE(120), QTTITLE2(120), QLIST(6,20), QSCHAM(200),
66		QFACTM(200), QFACT, QACT, QACT, QACT, QACT, QACT, QACT,
67	REAL	QFACTM, QFACT, QACT, QACT, QACT, QACT, QACT, QACT,
68	INTEGER	QCVERS, QLINE, QSLIME, QNSAS, QNSE, QKS, QNA, QNB, QKP, QNSS
69	INTEGER	QFRSTL, QLASTL, QSTEPL, QFRSTP, QLASTP, QSTEPP
70	INTEGER	QSTEP, QNCHAN, QNSS, QNCHAN, QNDACH, QFILE, QSPARE, QSCHAM
71	LOGICAL	QBNONE, QBPAST, QBPOLY, QBTASE, QBLSEG, QBLSEQ, QBLSEQ
72	LOGICAL	QFLDID, QTTITLE, QTTITLE2, QLIST
73	EQUIVALENCE	(QNSAS(1), QFRSTL, QNSA), (QNSAS(2), QLASTL, QNSB)
74	EQUIVALENCE	(QNSAS(3), QSTEPL, QKS), (QNSAS(4), QFRSTP, QNA)
75	EQUIVALENCE	(QNSAS(5), QLASTP, QNB), (QNSAS(6), QSTEPP, QKP)
76	LOGICAL	QCMOOO, QCM999
77	EQUIVALENCE	(QCMOOO, QCVERS)
78		
79		
80		
81		
82	C	COMMON /QANCL/ QANCLM, QANHEAD(200), QNCIL(180), QAN999
83	C	INTEGER QANCLM, QANHEAD, QNCIL
84	C	LOGICAL QANOOO, QAN999
85	C	EQUIVALENCE (QANOOO, QANCLM)
86		
87		
88		
89		
90	C	PARAMETERS* FOR U-2c IN FORTRAN PROGRAMS
91	INTEGER	QOCHAN, QOCHW2, QONSS, QONSS2, QOIBUF, QOIBUF
92	INTEGER	QOILM, QOMX2, QOMX4, QOLDTA, QOLDT2, QOLCTL
93	INTEGER	QUDATA, QURTRY, QUPRINT, QUERR, QUPNCH, QUPRINT, QURAP
94	INTEGER	QOVRTX, QONREG, QONSEG
95	INTEGER	QOVRER(4)
96	C	PARAMETERS --- DECLARATION SIZES FOR DATA ARRAYS
97	DATA	QOCHAN/200/
98	DATA	QOCHW2/24/
99	DATA	QONSS /3264/
100	DATA	QONSS2/824/
101	DATA	QOIBUF/22293/
102	DATA	QOIBUF/22293/
103	DATA	QOIBUF/22293/
104	DATA	QOVRTX/101/
105	DATA	QONREG/24/
106	DATA	QONSEG/24/
107	C	PARAMETERS --- MACHINE DEPENDENT LENGTH ATTRIBUTES, ETC.
108	DATA	QOILM/4/
109	DATA	QOMX2 /2/
110	DATA	QOMX4 /4/
111	DATA	QOLDTA/4/
112	DATA	QOLDT2/4/
113	DATA	QOLCTL/4/
114	C	PARAMETERS --- STANDARD I/O UNITS
115	DATA	QUDATA/4/
116	DATA	QURTRY/5/
117	DATA	QUPRINT/6/
118	DATA	QUERR /6/
119	DATA	QUPNCH/7/
120	DATA	QUPRINT/8/

```

DATA QUMAP /B/
C *PARAMETERS -- COMMON BLOCK VERSION TO CHECK THAT ALL MODULES COMPILED
C TO MATCH
C DATA QOCVER/'VNA3','MOO2','77/O','B/15'/
C -----
C -----
C -----
C -----
C INTEGER VERMES(14)
C
C WHEREVER "10" APPEARS IN THE FOLLOWING DECLARATIONS, IT CAN BE
C REPLACED BY THE NUMBER OF CHANNELS. WHEREVER "344" APPEARS, IT
C MEANS NUMBER OF POINTS.
C
C REAL L(MCHAN,MSS),MUO,MUOSO,MU,MU,MUSO,MFACT,EFACI(10,344),LPATH(1
C 10,344)
C REAL OM(10),ETA(10),A(10),B(10),SIG(10),LAMP(10),PHIP(10),
C 1 AT(10),BT(10),ATP(10),BTP(10),BTTP(10),CTPP(10),DTPP(10),
C 2 EO(10),E(10),FE(10),FP(10),TAU(10),RHO(10),ALPHA(10),
C 3 G(10),TAUZ(10),MUOQ(10),LBEAM(10),LSURF(10),FSCAT(10),
C 4 RHOS(10),LINTR(10),ATPP(10),D(10),WAVE(10),TAS(10)
C 5 TA(10),TR(10),TAUO3(10),TAUZO3(10),TAZ(10),TRZ(10)
C INTEGER PHID,PHIM,FSTP,LSTP,PTINC,SCATT
C INTEGER OPTION
C REAL FRES(344),LRSKY(10,344),LVIRP(10,344)
C REAL ED(10)
C REAL INDEX,THETA,REFRAC,AX,BX,SINA,SINB,TANA,TANB,FRESUN,FRESO,
C 1 EFACI,EFACTS,CONST,EXPO1,EXPO2,COND,COND1,COND2,COND3,
C 2 ARGSKY,ARGVIR,PEFSKY,PEVIR,MUFAC1,MUFAC2,ETOTF
C
C INDEX=REFRACTIVE INDEX OF THE LAKE WATER. WE SET INDEX=1.33. THE
C REFRACTIVE INDEX FOR PURE WATER, BUT YOU MAY WISH TO REPLACE
C THIS WITH EMPIRICAL VALUES.
C FRESO=THE LIMITING CASE OF THE FRESNEL REFLECTANCE AT ZERO ANGLE
C OF INCIDENCE. THIS WILL VARY WITH "INDEX."
C WHEN MU APPROACHES MUO OR TAU(1) APPROACHES TAUZ(1), THE STAN-
C DARD SINGLE SCATTERING RADIANCE FORMULAS TRUNCATE MOST OR ALL
C OF THE SIGNIFICANT FIGURES, AND THE EXPONENTIALS IN THEM MUST
C BE EXPANDED. "COND" CONTAINS THE VALUE TO CHECK AGAINST TO DE-
C CIDE WHEN TO USE THE EXPANSIONS. SET COND=10**(N-M) WHERE N=
C NUMBER OF SIGNIFICANT FIGURES DESIRED, M=MAXIMUM NUMBER OF FI-
C GURES HELD BY THE MACHINE. WE TAKE COND=0.01. IN REAL*4 PRE-
C CISION AND EXPANSIONS TO FOURTH ORDER TERMS, THIS WILL GUARAN-
C TEE RADIANCES TO AT LEAST THREE SIGNIFICANT FIGURES.
C
C DATA INDEX/1.33/,FRESO/0.0200593122/,COND/0.01/
C INTEGER CONTRL(MSS)
C DATA PRIOR/0./
C .....
C DATA VERMES/'ATMS','FR','TURN','ER A','TMO5','PHER','IC M','ODEL
C 1','RE','T/SH','W VO','0 6','/16','770'/
C DATA OML/.9999/
C DATA EPS/0.001/,PI/3.141592654/
C
C THESE "DATA" STATEMENTS ARE FOR DEBUGGING ONLY.
C
C DATA WAVE/ 428. 466. 508. 549. 592. 632. 674. 714. 756.
C 794/
C DATA LSW/6/
C DATA Z/12.323/

```

TABLE 2 (Cont.)

```

C DATA RHO/ .02 .02 .02 .02 .02 .02 .02 .02 .02 .02 /
C
C
C NOTE THAT WHEN YOU ENCOUNTER THE CALLS TO FUNCTION PF, VIZ
C PF(ARGUMENT,I,TR(I),TAS(I))
C THE FOURTH ARGUMENT SHOULD BE TAS(I), NOT TA(I), WHERE
C TAS(I)=FSCAT(I)*TA(I)
C SO THAT TAS(I) IS THE AEROSOL OPTICAL THICKNESS DUE ONLY TO SCAT-
C TERING.
C
C CON = 7.9577472E-02
C ATMSFR = 0
C CALL QOCALD(PRIOR,DATA,CONTRL,OMCHAN,OMSS)
C IF (QNONE OR QBPAS1)RETURN
C GO TO (1.99,3.99,99.6,99.9,99.99,99.99), QSTEP
C RETURN
C
C QSTEP-1 -- CONSTRUCT OLNE LINKAGE TO PRIOR MODULE SUPPLYING DATA:
C
C PRIOR = L(1,1)
C RETURN
C
C QSTEP-3 -- INITIALIZE; PRINT MODULE IDENTIFICATION, READ CONTROL
C DATA:
C
C CONTINUE
C CALL QOVERS TO PRINT MODULE NAME, DESCRIPTION, VERSION, DATE,
C PROGRAMMER,
C AND TO CHECK THAT THE COMMON BLOCK VERSION "QOCVER" IS CURRENT
C CALL QOVERS(VERMES,QOCVER)
C DON'T FORGET TO REMOVE THE "C" IN COLUMN 1 OF THE ABOVE COMMENT WHEN
C INTERFACING WITH "ATCOR."
C I HAVE DELETED THE CALL TO QOVERS IN ORDER TO RUN ATMSFR
C ALONE FOR DEBUGGING PURPOSES. JIM LIEZER
C WRITE (QUPRINT,100)
C 100 FORMAT('QATMSFR: ENTER FIRST PIXEL, LAST PIXEL, PIXEL INCREMENT,
C 1 SCATT, AND OPTION (515)')
C WRITE(QUPRINT,101)
C 101 FORMAT(' WHERE SCATT=0 FOR MULTISCATTERING,
C 1 SCATT=1 FOR SINGLE SCATTERING,
C 2 OPTION=0 FOR LPATH ONLY,
C 3 OPTION=1 FOR LPATH+LRSKY,
C 4 OPTION=2 FOR LPATH+LVIRP,
C 5 OPTION=3 FOR LPATH+LRSKY+LVIRP.')
C READ(QDATA,105)FSTP,LSTP,PTINC,SCATT,OPTION
C 105 FORMAT(515)
C RETURN
C
C QSTEP-6 -- BEGINNING OF REGION:
C
C CONTINUE
C
C COMPUTE WAVELENGTH DEPENDENT CONSTANTS
C
C CALL DATE(QDATA,QUPRINT,NDAY)
C CALL ANGLE(QDATA,QUPRIT,NDAY,PHIO,THETAO)
C READ HEIGHT OF SCANNER, HEIGHT OF GROUND, PRINT/WRITE TAPE SWITCH
C READ(QDATA,263) ZSCAN,ZGRND,LSW
C 263 FORMAT(2F8.5,15)
C ZSC=ZSCAN*3280.833

```

TABLE 2 (Cont.)

```

241 WRITE(QUPRNT,253)ZSC
242 FORMAT(' ',19X,'SENSOR ALTITUDE',F10.1,' FEET')
243 Z = ZSCAN-ZGRND
244 READ SCANNER WAVELENGTHS, ESTIMATED BACKGROUND REFLECTANCES
245 READ (QDATA,264) (WAVE(1),1-1,QNCHAN)
246 READ (QDATA,264) (RHO(1),1-1,QNCHAN)
247 FORMAT(10F8.5)
248 CALL PHASE(QDATA,QUPRNT,QNCHAN,WAVE)
249 CALL RAYLEI(QDATA,QUPRNT,QNCHAN,WAVE,TR,TAZ)
250 CALL OZONE(QDATA,QUPRNT,QNCHAN,WAVE,Z,TAUO3,TAUZO3)
251 CALL THICK(QDATA,QUPRNT,QNCHAN,WAVE,TAU)
252 CALL PARAMS(QDATA,QUPRNT,TR,TAUO3,TAU,QNCHAN,ETA,OM,TAS,FSCAT)
253 CALL AERO(QDATA,QUPRNT,QNCHAN,WAVE,Z,TAU,TR,TAUO3,TAZ,TA)
254 MUO = COS(THETAO)
255 CO2 = 2.*MUO
256 MUOSQ = MUO*MUO
257 CALL SOLAR(WAVE,QNCHAN,NDAY,EO)
258
259 C
260 C
261 C
262 C
263 C
264 READ THE AZIMUTH ANGLE COUNTERCLOCKWISE FROM NORTH TO ( THE FIRST
265 PIXEL OF THE ) SCAN PLANE.
266 READ (QDATA,265) PHID,PHIM,PHIS
267 WRITE (QUPRNT,270) PHID,PHIM,PHIS
270 FJRMAT('O',19X,'AZIMUTH ANGLE MEASURED COUNTERCLOCKWISE FROM NORTH',
271 'H TO FIRST PIXEL',/,',28X','OF SCAN PLANE IS',',13',
272 ' DEGREES',',13', MINUTES',',F5.1', SECONDS'.')
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265 FORMAT(213,F6.3)
266 PHI = 1.7453293E-02*PHID + 2.9088821E-04*PHIM + 4.8481368E-06*PHIS
267 PHI = SORT(1.0 - MUO*MUO)*COS(PHI-PHI0)
270 IF (SCATT.EQ.O.OR.OPTION.LT.2) GO TO 301
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1 'TIMER SQUARED'/' ' , .14X, 'PER MICROMETER PER STERADIAN.' /
2 11X, 'ALL IRRADIANCES ARE IN UNITS OF MILLIWATTS PER CENTIMETER'
3 ' ' , SQUARED'/15X, 'PER MICROMETER.' /
4 11X, 'ALL WAVELENGTHS ARE IN MICROMETERS.' )
IF (SCATT.EQ.O) GO TO 36
WRITE(QUPRNT,31)
31 FORMAT(' ',10X,'WE SHALL CALCULATE SINGLY SCATTERED RADIANCES ' ,
1 'ONLY.')
IF (OPTION.EQ.O) WRITE(QUPRNT,32)
32 FORMAT(' ',10X,'ONLY THE DIRECT PATH RADIANCE WILL BE CONSIDERED.
1 ')
IF (OPTION.EQ.1) WRITE(QUPRNT,33)
33 FORMAT(' ',10X,'THE DIRECT PATH RADIANCE AND THE REFLECTED SKY ' ,
1 'RADIANCE WILL BE CONSIDERED.')
IF (OPTION.EQ.2) WRITE(QUPRNT,34)
34 FORMAT(' ',10X,'THE DIRECT PATH RADIANCE AND THE VIRTUAL PATH ' ,
1 'RADIANCE WILL BE CONSIDERED.')
IF (OPTION.EQ.3) WRITE(QUPRNT,35)
35 FORMAT(' ',10X,'THE DIRECT PATH RADIANCE AND BOTH THE VIRTUAL ' ,
1 'PATH RADIANCE AND THE REFLECTED SKY RADIANCE ' ,
2 'WILL BE CONSIDERED.')
GO TO 38
36 WRITE(QUPRNT,37)
37 FORMAT(' ',10X,'WE SHALL CALCULATE MULTIPLY SCATTERED RADIANCES ' ,
1 'ONLY.')
38 CONTINUE
WRITE(QUPRNT,275) OBGANG,QDANG
275 FORMAT(20X,'MAXIMUM SCAN ANGLE=' ,F9.7,' RADIANCS.
1 SOLUTION OF SCANNER IS ' ,F9.7,' RADIANCS.')
WRITE(QUPRNT,276)(I=1,ONCHAN)
276 FORMAT('1 CHANNEL ' ,10I10)
WRITE(QUPRNT,277)(WAVE(I),I=1,ONCHAN)
277 FORMAT(' - WAVELENGTH ' ,10F10.5)
WRITE(QUPRNT,278)(TAU(I),I=1,ONCHAN)
278 FORMAT(' - INTERPOLATED OPTICAL ' ,10F10.5)
WRITE(QUPRNT,280)
280 FORMAT(' THICKNESS')
WRITE(QUPRNT,281)(RHO(I),I=1,ONCHAN)
281 FORMAT(' - BACKGROUND ALBEDO ' ,10F10.5)
WRITE(QUPRNT,283)(OM(I),I=1,ONCHAN)
283 FORMAT(' - SINGLE SCATTERING ' ,10F10.5)
WRITE(QUPRNT,284)
284 FORMAT(' ALBEDO')
WRITE(QUPRNT,285)(FSCAT(I),I=1,ONCHAN)
285 FORMAT(' - SCATTERING PARAMETER ' ,10F10.5)
WRITE(QUPRNT,287)
287 FORMAT(' (FSCAT)')
WRITE(QUPRNT,290)(TAUO3(I),I=1,ONCHAN)
290 FORMAT(' - OZONE OPTICAL ' ,7X,10F10.5)
WRITE(QUPRNT,292)
292 FORMAT(' THICKNESS')
WRITE(QUPRNT,295)(TR(I),I=1,ONCHAN)
295 FORMAT(' - RAYLEIGH OPTICAL ' ,4X,10F10.5)
WRITE(QUPRNT,296)
296 FORMAT(' THICKNESS')
DO 10 I=1,ONCHAN
TAUZ(I) = TAUZ03(I) + TRZ(I) + TAZ(I)
IF (SCATT.EQ.O) GO TO 500

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COMPUTE WAVELENGTH DEPENDENT QUANTITIES NEEDED FOR

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TABLE 2 (Cont.)

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SINGLE SCATTERING CALCULATIONS
SIG(I) = TAU(I) - TAUZ(I)
EFACIO = EXP(SIG(I)/MUQ)
CONST = CON*OM(I)*EO(I)
EFACIS = EXP(-TAU(I)/MUQ)
ETOTF = (1.-ETA(I))*TAU(I)
E(I) = (MUQ*MUO*EO(I)/(MUO*ETOTF))*(1+(2.*RHO(I)*ETOTF)/
(1+2.*ETOTF))

FE(I)=MUO*EO(I)/E(I)
COMPUTE DIRECT IRRADIANCE AT SURFACE
ED(I)=MUO*EO(I)*EFACIS
GO TO 15
CONTINUE

500

COMPUTE WAVELENGTH DEPENDENT QUANTITIES NEEDED FOR
MULTI SCATTERING CALCULATIONS

IF (OM(I).GT. OML) GO TO 20
A(I) = OM(I)* (1.0-ETA(I))
B(I) = 1.0 + A(I) - OM(I)
C = A(I) + B(I)
NU = MUO / SORT(C*(B(I)-A(I)))
NUSQ(I) = NU*NU
SIG(I) = TAU(I) - TAUZ(I)
ARG1 = SIG(I)/NU
SH1 = SINH(ARG1)
CH1 = COSH(ARG1)
ARG2 = TAU(I)/NU
SH2 = SINH(ARG2)
CH2 = COSH(ARG2)
ARG3 = C02*ARG2
SH3 = SINH(ARG3)
CH3 = COSH(ARG3)
ARG4 = ARG3*TAUZ(I)/TAU(I)
SH4 = SINH(ARG4)
CH4 = COSH(ARG4)

COMPUTE LAMBDA DOUBLE PRIME
LAMPPI(I) = CON*NUSQ(I) * OM(I)*EO(I) / (B(I)*NU*SH2 + MUO*CH2)

COMPUTE PHI PRIME
PHIP(I) = 2.0*RHO(I)*MUOSQ*LAMPPI(I)/
(MUO*CH3+NU*(B(I)-A(I)*RHO(I))*SH3)

COMPUTE OTHER CONSTANTS
AT(I) = NU*SH1
BT(I) = CH1
ATP(I) = B(I)*NU*SH1 + MUO*CH1
BTP(I) = B(I)*CH1 + MUO*SH1/NU
ATPP(I) = C*NU*SH4 + MUO*CH4
BTPP(I) = 2.0*C*((1.0-OM(I))*NU*SH4 + MUO*CH4)
CTPP(I) = C*NU*SH3 + MUO*CH3
DTPP(I) = 2.0*C*((1.0-OM(I))*NU*SH3 + MUO*CH3)

COMPUTE TOTAL IRRADIANCE AT SURFACE
E(I) = 6.28318531*PHIP(I) + (NU*B(I)*SH3 + MUO*CH3)/
(NUSQ(I)*OM(I)*RHO(I))
FE(I) = MUO*EO(I)/E(I)

COMPUTE DIRECT IRRADIANCE AT SURFACE
ED(I)=MUO*EO(I)*EXP(-TAU(I)/MUO)
GO TO 15

CALCULATE SIMILAR CONSTANTS FOR OMEGA=1
CONTINUE

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TABLE 2 (Cont.)

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421 SF = 1.0-ETA(I)
422 D(I) = MUO/SF
423 FP(I) = CON*SF*EO(I) / (MUO + SF*TAU(I))
424 CHI = CO2*MUO*RHO(I) / (1.0 + 2.0*SF*(1.0-RHO(I))*TAU(I))
425 ALPHA(I) = CHI* (4.0 + 1.0/SF)
426 G(I) = 4.0*CHI
427 SIG(I) = TAU(I) - TAUZ(I)
428 COMPUTE TOTAL IRRADIANCE AT SURFACE
429 E(I) = 6.28318531*CHI*FP(I)*(1.0 + 2.0*SF*TAU(I))/(RHO(I)*SF)
430 FE(I) = MUO*EO(I)/E(I)
431 COMPUTE DIRECT IRRADIANCE AT SURFACE
432 ED(I) = MUO*EO(I)*EXP(-TAU(I)/MUO)
433 CONTINUE
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TABLE 2 (Cont.)

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181 IF (THETA) 58, 59, 58
182 THETA=ABS(THETA)
183 REFRACT=ARSIN(SIN(THETA))*1./INDEX
184 AX=THETA-REFRACT
185 BX=THETA+REFRACT
186 SINA= SIN(AX)
187 SINB= SIN(BX)
188 TANA= TAN(AX)
189 TANB= TAN(BX)
190 FRES(IP)=0.5*((SINA*SINA)/(SINB*SINB)+(TANA*TANA)/(TANB*TANB))
191 GO TO 310
192 FRES(IP)=FRESO
193 CONTINUE
194
195 C
196 C
197 C
198 C
199 C
200 ARGSKY=FMU-FPHI*ROOT
201 PFISKY=PF(ARGSKY,I,TR(I),TAS(I))
202 EXP02=TAU(I)/MUJ
203 MUJFAC2=(MUJ-MU)/MUJ
204 COND2=EXP02*MUJFAC2
205 IF (A-.(COND2).LT.COND) GO TO 503
206
207 C
208 C
209 C
210 C
211 C
212 CALCULATE THE REFLECTED SKY RADIANCE
213
214 LRSKY(I,IP)=FRES(IP)*(CONST/MUJFAC2)*PFISKY*EFACT(I,IP)*
215 (EFACTS-EXP(-EXP02))
216 GO TO 505
217
218 C
219 C
220 C
221 C
222 C
223 EITHER MU AND MUO OR TAU(I) AND TAUZ(I) ARE TOO CLOSE TOGETHER
224 SO THE FOURTH ORDER EXPANSION OF THE REFLECTED SKY RADIANCE
225 MUST BE USED.
226
227 LRSKY(I,IP)=FRES(IP)*CONST*PFISKY*EFACT(I,IP)*EFACTS*EXP02*(1.
228 -(COND2/2.)*(1.-(COND2/3.))*(1.-(COND2/4.))
229 )
230 IF (OPTION.LT.2) GO TO 2019
231
232 C
233 C
234 C
235 C
236 C
237 CALCULATE THE SCAN ANGLE DEPENDENT QUANTITIES NEEDED TO FIND
238 THE PATH RADIANCE GENERATED BY THE VIRTUAL SUN.
239
240 ARGVIR=FMU-FPHI*ROOT
241 PFVIR=PF(ARGVIR,I,TR(I),TAS(I))
242 MUJFAC2=(MUJ-MU)/MUJ
243 COND3=EXP01*MUJFAC2
244 IF (ABS(COND3).LT.COND) GO TO 506
245
246 C
247 C
248 C
249 C
250 C
251 CALCULATE THE PATH RADIANCE GENERATED BY THE VIRTUAL SUN.
252
253 LV[RP(I,IP)]=FRESUN*(CONST/MUJFAC2)*PFVIR*EFACTS*(1./EFACTO -
254 EFACT(I,IP))
255 GO TO 2019
256
257 C
258 C
259 C
260 C
261 C
262 EITHER MU AND MUO OR TAU(I) AND TAUZ(I) ARE TOO CLOSE TOGETHER
263 SO WE USE THE EXPANSION OF THE VIRTUAL PATH RADIANCE.
264
265 LVIRP(I,IP)=FRESUN*CONST*PFVIR*EFACTS*EFACT(I,IP)*EXP01*(1.+
266 (COND3/2.)*(1.+(COND3/3.)*(1.+(COND3/4.))))
267 GO TO 2019
268
269 C
270 C
271 C
272 C
273 C

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TABLE 2 (Cont.)

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541      510      CONTINUE
542      C
543      C      CALCULATE THE SCAN ANGLE DEPENDENT QUANTITIES NEEDED TO FIND
544      C      THE MULTIPLY SCATTERED RADIANCE.
545      C
546      IF (OM(I) .GT. OML) GO TO 52
547      C
548      C      DEFINE SINGULARITIES
549      SING1 = NUSQ(I) - MUJQ
550      SING2 = NUSQ(I) - CO2*CO2*MUSQ
551      IF (ABS(SING1) .GT. EPS .AND. ABS(SING2) .GT. EPS) GO TO 55
552      WRITE (QUERR,380)
553      FORMAT(' -ATMSFR: *****ERROR-- SINGULARITY EXISTS IN EQUATI
554      ONS*****')
555      ATMET = 57.29578*THETA
556      WRITE (QUERR,385) ATMET
557      FORMAT(' -ATMSFR: SINGULAR AT THETA= ',F7.4)
558      GO TO 2019
559      55      CONTINUE
560      PF1 = PF(ARG,I,TR(I), TAS(I))
561      PF2 = PF(-ARG,I,TR(I), TAS(I))
562      EFAC(I,IP) = EXP(-SIG(I)/MU)
563      MFAC = MU * EFAC(I,IP)
564      C      CALCULATE PATH RADIANCE
565      LPATH(I,IP) = LAMPP(I) *
566      (A(I)*(AT(I)-BT(I)*MU + MFAC) *PF1
567      + (ATP(I) -BTP(I)*MU -MUJQ*EFAC(I,IP) +B(I)*MFAC)*PF2)/SING1
568      +PHIP(I)*
569      (ATPP(I) +BTTP(I)*MU -CTPP(I)*EFAC(I,IP) -DTPP(I)*MFAC)
570      /SING2
571      GO TO 2019
572      52      CONTINUE
573      C      CALCULATE PHASE FUNCTIONS
574      PF1 = PF(ARG,I,TR(I), TAS(I))
575      PF2 = PF(-ARG,I,TR(I), TAS(I))
576      EFAC(I,IP) = EXP(-SIG(I)/MU)
577      CFAC = 1.0 -EFAC(I,IP)
578      C      CALCULATE PATH RADIANCE
579      LPATH(I,IP) = FP(I) *
580      ((SIG(I) -MU*CFAC)*PF1 + (SIG(I) + (D(I)-MU)*CFAC)*PF
581      2 + (ALPHA(I) + G(I)*MU)*CFAC)
582      2019      CONTINUE
583      10      CONTINUE
584      C
585      C      ECHO MORE QUANTITIES FOR THE USER'S CONVENIENCE.
586      C
587      WRITE (QUPRNT,555)(ETA(I),I=1,QNCHAN)
588      FORMAT(' - ANISOTROPY PARAMETER ',10F10.5)
589      WRITE (QUPRNT,291)
590      (ETA)
591      WRITE (QUPRNT,282)(TAUZ(I),I=1,QNCHAN)
592      FORMAT(' - OPTICAL DEPTH ',10F10.5)
593      WRITE (QUPRNT,297)(E(I),I=1,QNCHAN)
594      FORMAT(' - TOTAL IRRADIANCE ',4X,10F10.5)
595      WRITE (QUPRNT,298)
596      (ON SURFACE)
597      WRITE (QUPRNT,370)(ED(I),I=1,QNCHAN)
598      FORMAT(' - DIRECT IRRADIANCE ',10F10.5)
599      WRITE (QUPRNT,371)
600      (LN SURFACE)
601      RETURN

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TABLE 2 (Cont.)

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601 .....
602 C QSTEP=8 -- NEXT LINE TO BE PROCESSED: .....
603 .....
604 B CONTINUE .....
605 C IF WE ARE PRINTING RADIOMETRIC QUANTITIES, START EACH LINE ON .....
606 C A NEW PAGE .....
607 IF (LSW.EQ.6) WRITE(QUPRNT,400) .....
608 FORMAT('1') .....
609 DO 1009 IP=FSIP,LSIP,PTINC .....
610 LOOP FOR EACH PIXEL--QNSS=NUMBER OF PIXELS .....
611 C TEST "BAD DATA" FLAG FOR THIS PIXEL. IF NON-ZERO, DON'T PROCESS. .....
612 IF (CONTRL(IP).NE.0) GO TO 2009 .....
613 DO 3009 IW=1,QNCHAN .....
614 C TOTAL RADIANCE (EXPERIMENTAL) = L(IW,IP) .....
615 LBEAM(IW) = L(IW,IP) - LPATH(IW,IP) .....
616 IF (SCATT.EQ.0) GO TO 520 .....
617 IF (OPTION.EQ.1 OR OPTION.EQ.3) LBEAM(IW)=LBEAM(IW)-LRSKY(IW, .....
618 IP) .....
619 IF (OPTION.EQ.2) LBEAM(IW)=LBEAM(IW)-LVIRP(IW,IP) .....
620 CONTINUE .....
621 IF (L(IW,IP).EQ.0) LBEAM(IW) = 0 .....
622 CALCULATE SURFACE RADIANCE .....
623 LSURF(IW) = LBEAM(IW)/EFAC(IW,IP) .....
624 CALCULATE INTRINSIC RADIANCE .....
625 LINTR(IW) = FE(IW)*LSURF(IW) .....
626 CALCULATE SURFACE REFLECTANCE .....
627 RHOS(IW) = LSURF(IW)/E(IW) .....
628 GO TO (71,72,73,74,75,76).LSW .....
629 71 L(IW,IP) = LPATH(IW,IP) .....
630 GO TO 60 .....
631 72 L(IW,IP) = LBEAM(IW) .....
632 GO TO 60 .....
633 73 L(IW,IP) = LSURF(IW) .....
634 GO TO 60 .....
635 74 L(IW,IP) = LINTR(IW) .....
636 GO TO 60 .....
637 75 L(IW,IP) = RHOS(IW) .....
638 GO TO 60 .....
639 CONTINUE .....
640 THETA = (QBANG-QDANG*(IP-1)) *180/PI .....
641 CONTINUE .....
642 CONTINUE .....
643 IF (LSW.NE.6) GO TO 3019 .....
644 PRINT RADIOMETRIC QUANTITIES .....
645 IF (SCATT.EQ.0 OR OPTION.EQ.0 OR OPTION.EQ.2) WRITE(QUPRNT,410 .....
646 .....
647 FORMAT(5X,'LINE NUMBER',16.5X,'PIXEL NUMBER',14.5X, .....
648 'JOLINE,IP,THETA .....
649 'SCAN ANGLE-',F7.3,'DEGREES') .....
650 IF (SCATT.NE.0 AND (OPTION.EQ.411) JOLINE,IP,THETA,FRES(IP), .....
651 FORMAT(5X,'LINE NUMBER',16.5X,'PIXEL NUMBER',14.5X, .....
652 'SCAN ANGLE-',F7.3,'DEGREES .....
653 'FRESNEL REFLECTANCE-', .....
654 .2F6.3) .....
655 WRITE (QUPRNT,420) .....
656 FORMAT('O',14X,'1',9X,'2',9X,'3',9X,'4',9X,'5',9X,'6', .....
657 9X,'7',9X,'8',9X,'9',9X,'10') .....
658 WRITE (QUPRNT,425) (L(IW,IP),IW=1,QNCHAN) .....
659 FORMAT('O LTOT',10F10.5) .....
660 WRITE (QUPRNT,430) (LPATH(IW,IP),IW=1,QNCHAN) .....
661 FORMAT(' LPATH',10F10.5)

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TABLE 2 (Cont.)

561	IF (SCATT.EQ.0)GO TO 433
562	IF (OPTION.EQ.1 OR OPTION.EQ.3)WRITE(QUPRINT,431)(LRISKY(IW,IP)
563	,IW-1,ONCHAN)
564	FORMAT(' LRISKY ',10F10.5)
565	IF (OPTION.EQ.2)WRITE(QUPRINT,432)(LVIRP(IW,IP),IW-1,ONCHAN)
566	FORMAT(' LVIRP ',10F10.5)
567	WRITE(QUPRINT,434)(EFACT(IW,IP),IW-1,ONCHAN)
568	FORMAT(' TRANS ',10F10.5)
569	WRITE (QUPRINT,435) LBEAM
570	FORMAT(' LBEAM ',10F10.5)
571	WRITE (QUPRINT,440) LSURF
572	FORMAT(' LSURF ',10F10.5)
573	WRITE (QUPRINT,445) LINTR
574	FORMAT(' LINTR ',10F10.5)
575	WRITE (QUPRINT,450) RHOS
576	FORMAT(' RHOS ',10F10.5/'O')
577	SET QBNONE TRUE TO TELL ANY FOLLOWING MODULES THAT NO OUTPUT
578	DATA ARE PRESENT.
579	QBNONE = .TRUE.
580	3019 CONTINUE
581	2009 CONTINUE
582	1009 CONTINUE
583	RETURN
584	END

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The user must also define the center wavelengths (μm) of the multispectral scanner and the corresponding values of the surface background albedo (values from zero to one). Also, the center wavelengths (μm) of the surface radiometers and the corresponding optical thicknesses must be known. It should be noted that the optical thicknesses used should be those measured as closely as possible in time with the multispectral data.

5.4 MODEL CALCULATIONS

In this section we present several examples of the radiances for the various components. Because our main interest is in the radiance components as a function of scan angle and visibility, we will present the results of the calculations in terms of these parameters.

Figure 10 depicts the variation in the singly-scattered reflected sky radiance at the sensor as a function of the nadir scan angle and visibility. In this example the solar zenith angle is 45° and the scan plane is perpendicular to the solar plane. The curves which result are a combination of the variation of the sky radiance, the transmittance from the surface to the sensor, and the Fresnel reflectance of the water surface. For a practical scanner with a maximum scan angle of about 45° the curves indicate that one would not observe the large radiance peaks at the large angles.

In Figure 11 we display the corresponding path radiance as a result of singly-scattered radiation from the reflection of the sun in the water. In this case the radiance peaks do not exist at the large scan angles.

In Figure 12 we illustrate the relative magnitudes of the various radiation components as a function of scan angle for a moderately hazy atmosphere. The virtual sun path radiance is the smallest value and the multiply scattered sky radiance is the largest value.

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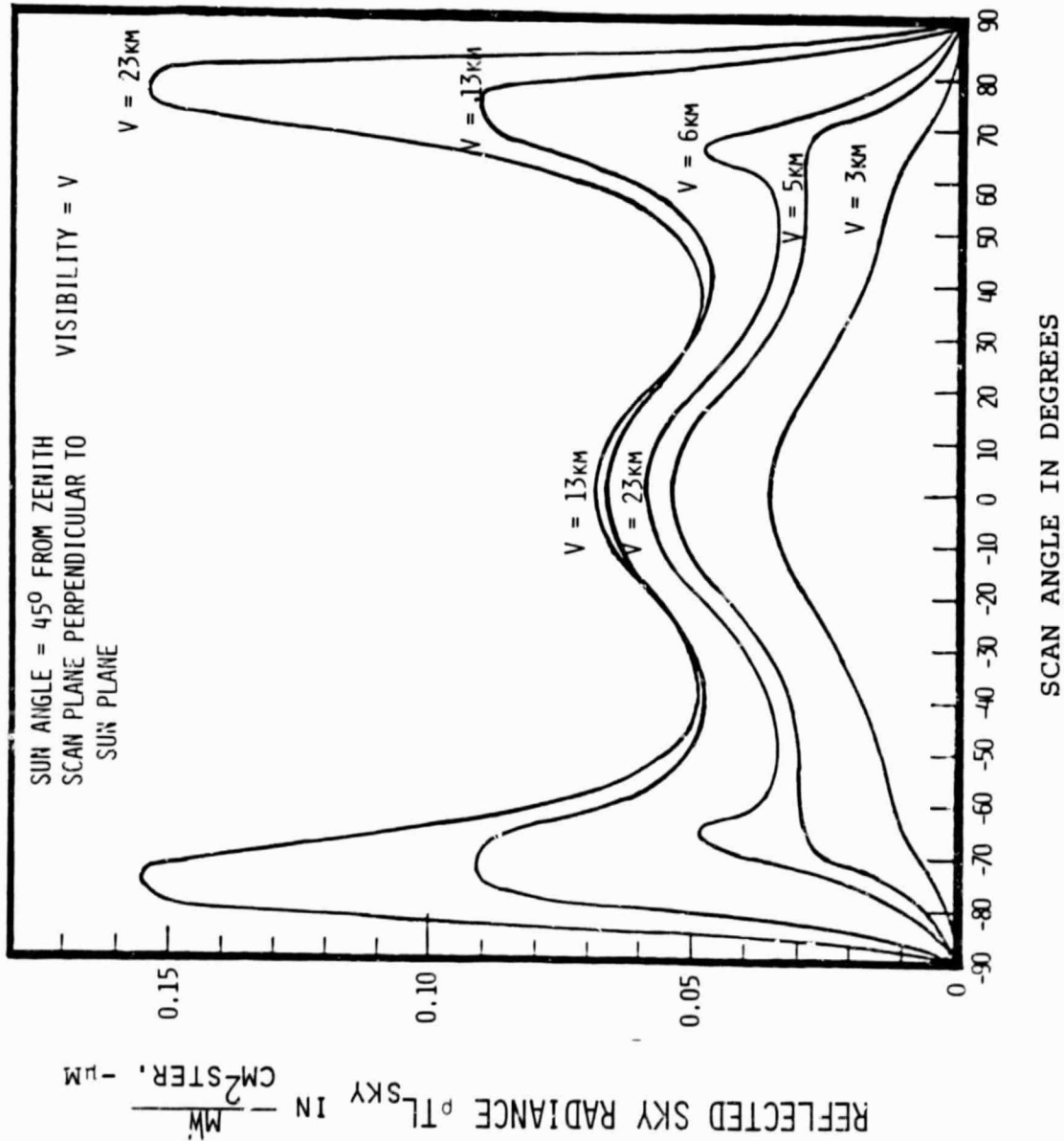


Figure 10. Singly-Scattered Reflected Sky Radiance at Sensor
for a Wavelength of $0.55 \mu m$

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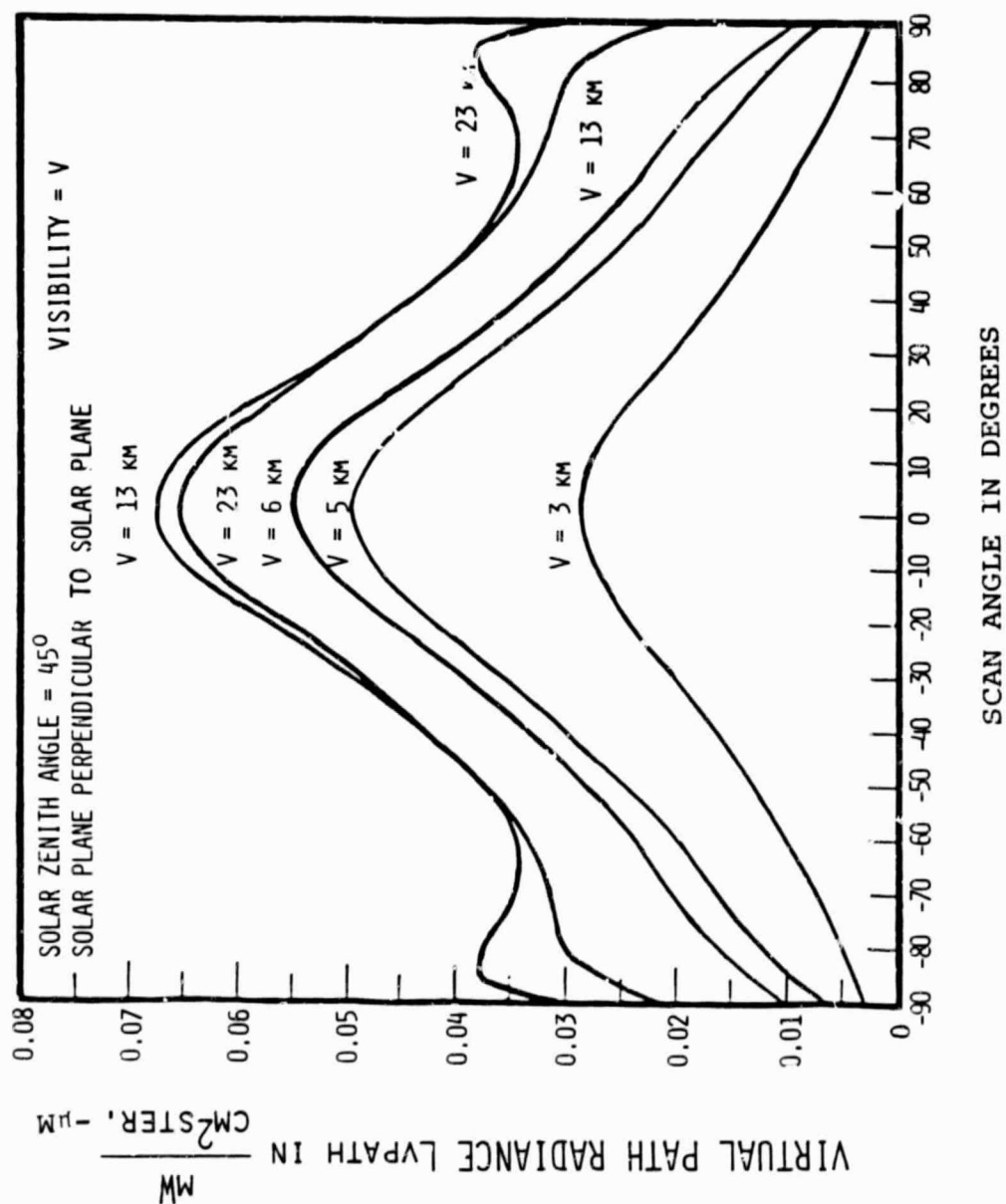


Figure 11. Singly-Scattered Virtual Path Radiance at Sensor
for a Wavelength of $0.55\mu\text{M}$

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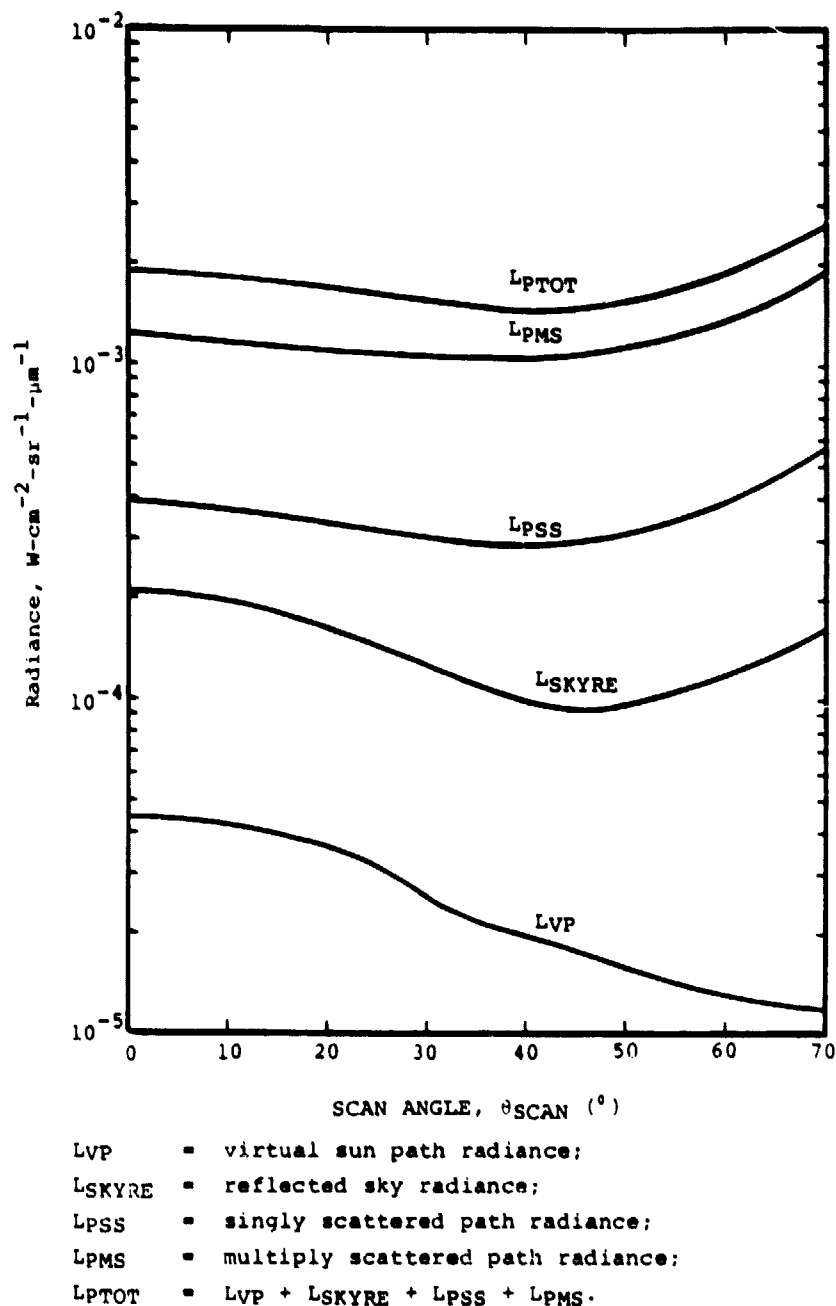


Figure 12. SURFACE AND PATH RADIANCE COMPONENTS DETECTED BY SENSOR AT OPTICAL DEPTH OF $\tau = 0.346$ AS A FUNCTION OF SCAN ANGLE, θ_{SCAN} . SCAN PLANE \perp SOLAR PLANE, $\theta_{\text{SUN}} = 30^{\circ}$, $\lambda = 0.55 \mu\text{m}$, PHASE FUNCTION = CONTINENTAL REFRACTIVE INDEX $1.5 - 0.01i$, VISIBILITY = 10 KM.

In Figure 13 we indicate the variation in the ratio of the singly-scattered sky radiance to the singly-scattered path radiance as a function of the optical depth τ of the sensor. As the curves illustrate, the reflected sky radiance component is relatively more important for the larger optical depths.

Figure 14 illustrates the variation in the ratio of the virtual sun path radiance to the singly-scattered path radiance with scan angle for four optical depths.

Figure 15 depicts the large ratio of the multiply-scattered component to the singly-scattered path radiance component as a function of optical depth and scan angle.

Because optical thickness or visibility is of major importance in remote sensing investigations, we want to consider the variation of the radiance components with respect to visibility. This effect is illustrated in Figure 16 for three different atmospheres. We chose the continental aerosol because it more nearly represents the type which would be found over the Great Lakes. The three refractive indices are: $1.5-0.0i$ which corresponds to a "clean" haze, i.e., one where there is no absorption; $1.5-0.01i$ which corresponds to a haze with some aerosol absorption; and $1.5-0.1i$, a complex index of refraction which corresponds to a haze with more absorption. As the curves indicate, an absorbing haze or one which corresponds to considerable air pollution gives rise to a large ratio of reflected sky radiance relative to the singly-scattered path radiance.

The effect of the complex index of refraction is also evident in the ratio of the virtual sun path radiance to the singly-scattered path radiance as indicated in Figure 17.

Finally, we illustrate in Figure 18 the variation of various combinations of ratios in terms of the visibility for a refractive index of $1.5-0.01i$.

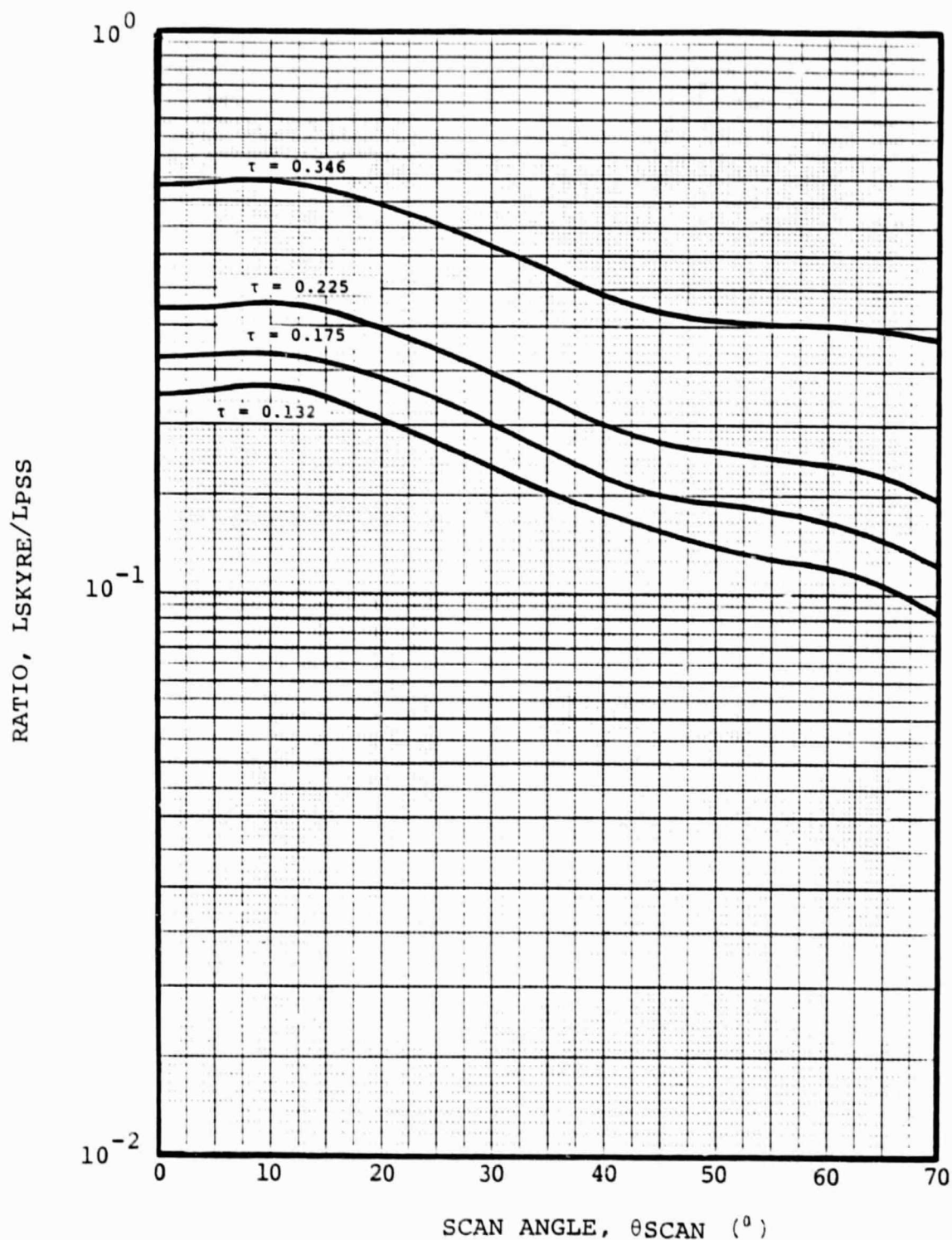


FIGURE 13. RATIO OF REFLECTED SINGLY SCATTERED SKY RADIANCE TO SINGLY SCATTERED PATH RADIANCE AS A FUNCTION OF SCAN ANGLE, θ_{SCAN} , FOR OPTICAL DEPTH, τ , OF THE SENSOR OF 0.132, 0.175, 0.225 and 0.346. SCAN PLANE \perp SOLAR PLANE, PHASE FUNCTION = CONTINENTAL REFRACTIVE INDEX $1.5 - 0.01i$, VISIBILITY = 10 KM, $\lambda = 0.55 \mu M$, $\theta_{SUN} = 30^\circ$.

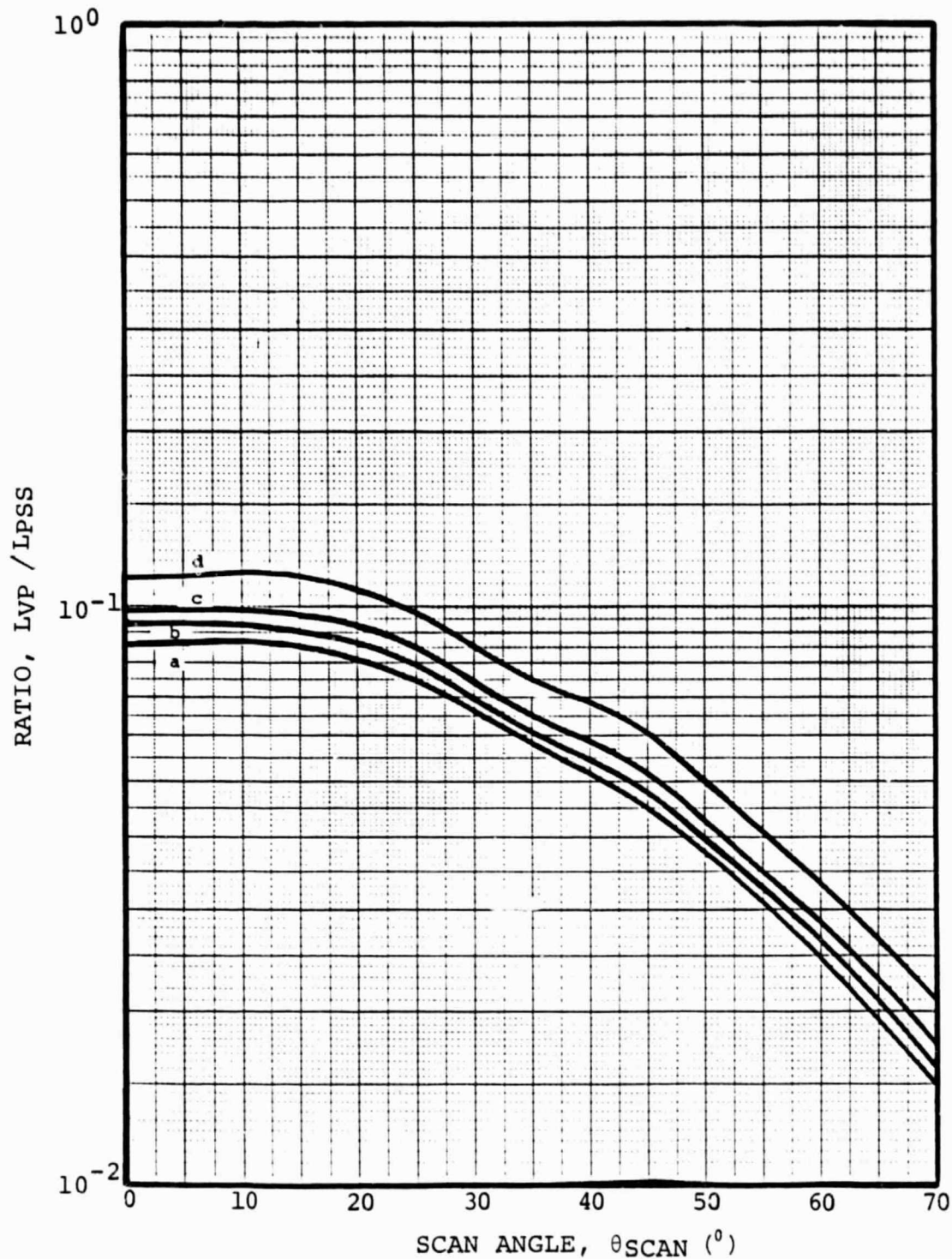


FIGURE 14. RATIO OF SINGLY SCATTERED PATH RADIANCE FROM THE VIRTUAL SUN, L_{VP} , TO SINGLY SCATTERED PATH RADIANCE, L_{PSS} , AS A FUNCTION OF SCAN ANGLE, θ_{SCAN} , FOR OPTICAL DEPTHS, τ , OF THE SENSOR OF A) 0.132, B) 0.175, C) 0.225 AND D) 0.346. SCAN PLANE \perp SOLAR PLANE, PHASE FUNCTION = CONTINENTAL REFRACTIVE INDEX $1.5 - 0.01i$, VISIBILITY = 10 KM, $\lambda = 0.55 \mu M$, $\theta_{SUN} = 30^\circ$.

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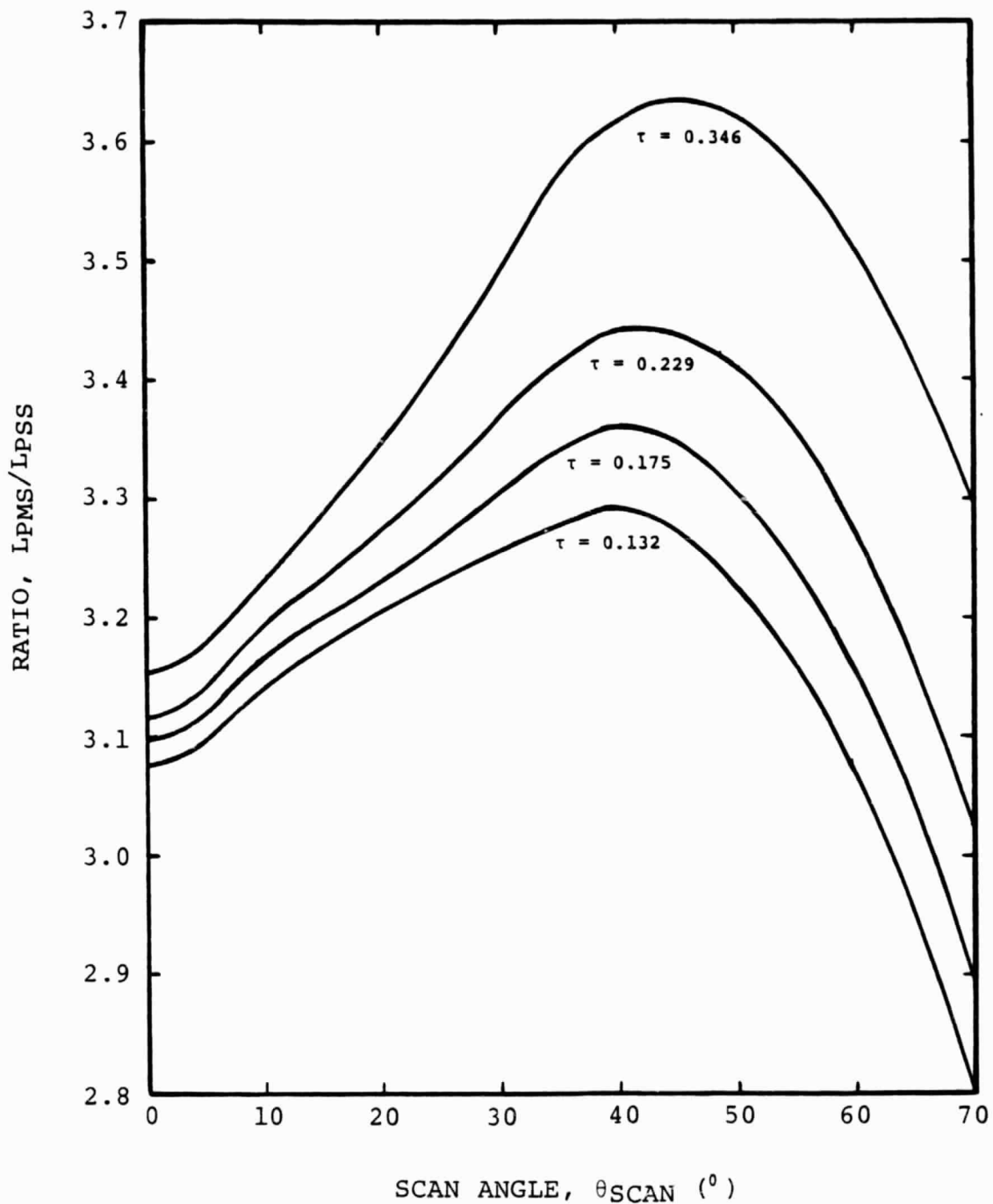


FIGURE 15. RATIO OF MULTIPLY SCATTERED PATH RADIANCE, LPMS, TO SINGLY SCATTERED PATH RADIANCE, LPSS, AS A FUNCTION OF SCAN ANGLE, θ_{SCAN} , FOR OPTICAL DEPTHS, τ , OF THE SENSOR OF 0.132, 0.175, 0.229, 0.346. SCAN PLANE \perp SOLAR PLANE, VISIBILITY = 10 KM, PHASE FUNCTION = CONTINENTAL REFRACTIVE INDEX $1.5 - 0.01i$, $\lambda = 0.55 \mu\text{M}$, $\theta_{\text{SUN}} = 30^\circ$.

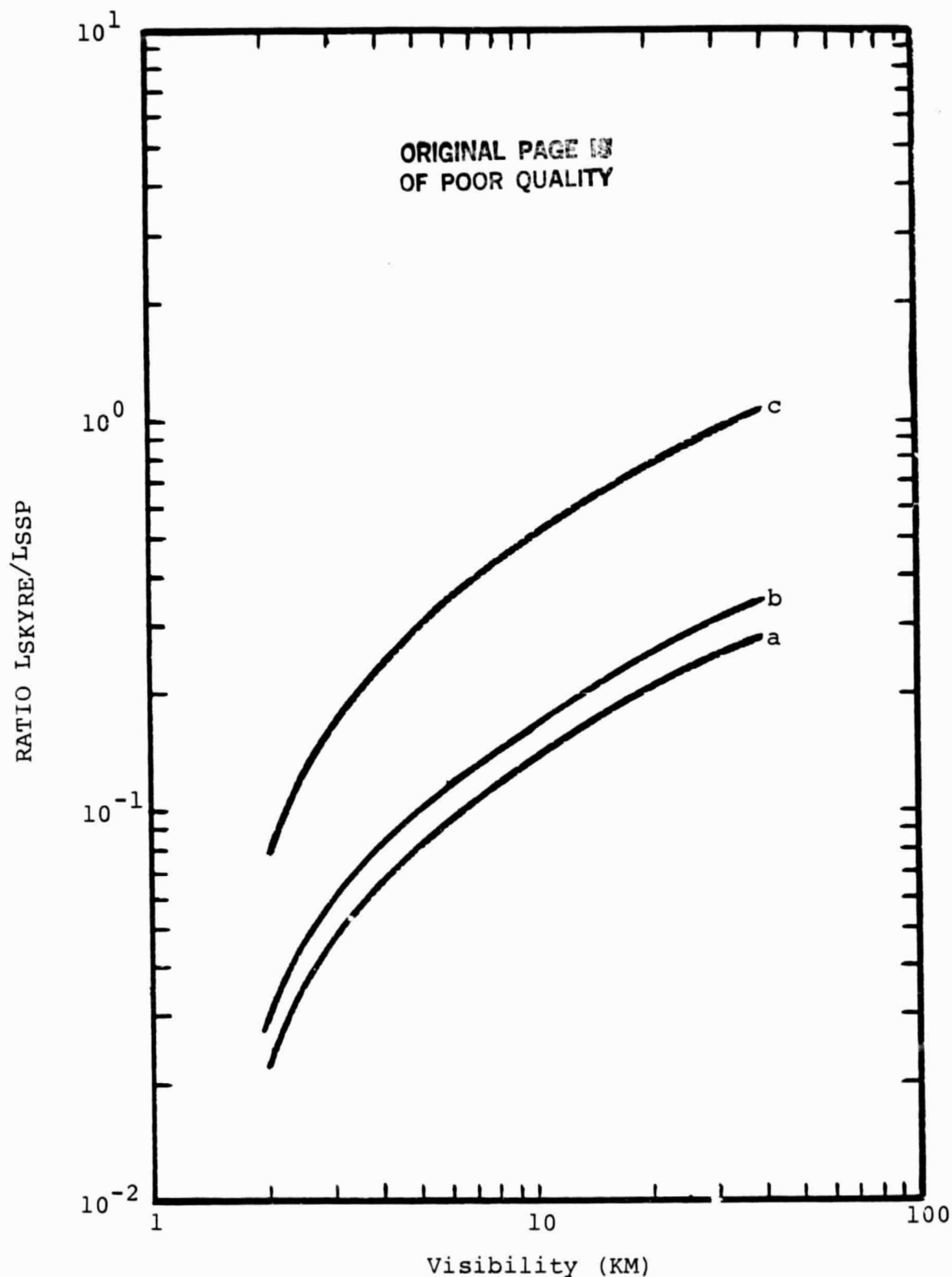


Figure 16. RATIO OF REFLECTED SKY RADIANCE TO SINGLY SCATTERED PATH RADIANCE FOR THREE CONTINENTAL AEROSOL MODELS: a) 1.5 - 0.0i, b) 1.5 - 0.01i, and c) 1.5 - 0.1i AS A FUNCTION OF ATMOSPHERIC VISIBILITY. SCAN PLANE \perp SOLAR PLANE, $\theta_{\text{SUN}} = 30^\circ$, $\lambda = 0.55 \mu\text{M}$.

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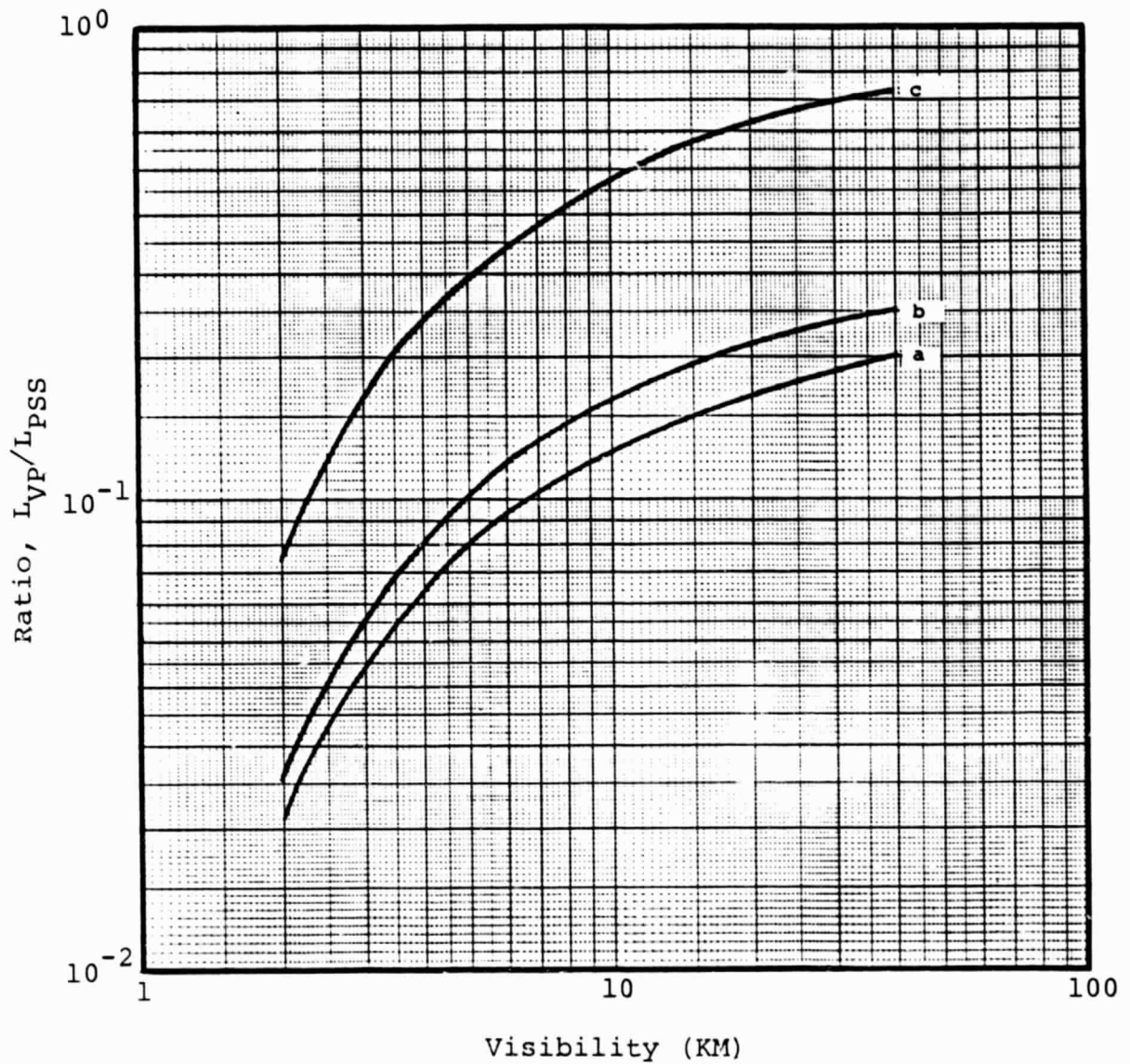
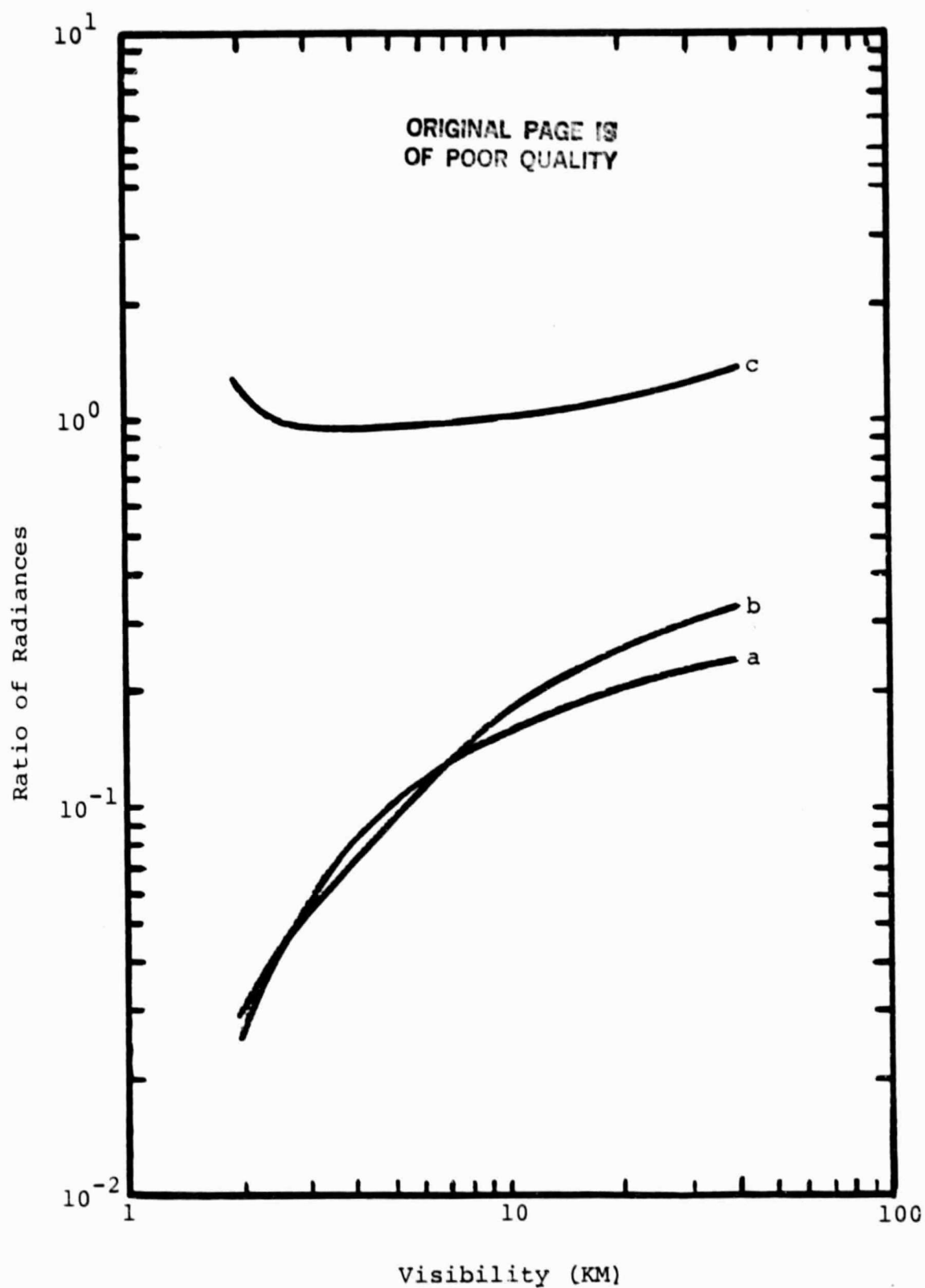


FIGURE 17. RATIO OF SINGLY SCATTERED PATH RADIANCE FROM THE VIRTUAL SUN, L_{VP} , TO SINGLY SCATTERED PATH RADIANCE, L_{PSS} , AS A FUNCTION OF ATMOSPHERIC VISIBILITY FOR THREE CONTINENTAL AEROSOL MODELS WITH REFRACTIVE INDICES OF a) $1.5 - 0.01i$, b) $1.5 - 0.01i$, and c) $1.5 - 0.1i$. $\lambda = 0.55 \mu\text{M}$, $\theta_{\text{SUN}} = 30^\circ$, $\theta_{\text{SCAN}} = 0^\circ$.



L_{vp} = virtual sun path radiance;
 L_{pss} = singly scattered path radiance;
 L_{skyre} = reflected sky radiance.

FIGURE 18. RATIOS OF (a) L_{vp}/L_{pss} , (b) L_{skyre}/L_{pss} , (c) L_{skyre}/L_{vp} AS A FUNCTION OF VISIBILITY (KM). $\lambda = 0.55 \mu m$, $\theta_{SUN} = 30^\circ$, $\theta_{SCAN} = 0^\circ$, AEROSOL MODEL = CONTINENTAL REFRACTIVE INDEX $1.5 - 0.01i$.

CONCLUSIONS AND RECOMMENDATIONS

The problem of developing an atmospheric correction algorithm for remote sensing is an old and difficult one. The main difficulty lies in not being able to have available sufficient data which can be used to specify the values of the relevant atmospheric parameters. The problem is all the more difficult in the case of the remote sensing of water bodies because of the low signal-to-noise ratio involved. In this investigation we have extended an existing computer algorithm so as to include additional radiation components. The original algorithm included the path radiance which arises from the singly-scattered solar radiation in the atmosphere. We have now included the radiance which arises from the sky radiation which is reflected by the water surface and is then attenuated as it propagates from the surface to the sensor. In addition, we have included the path radiance component which arises from the single scattering of radiation as a result of a virtual sun, i.e., of the sun's reflection in the water. It should be realized that this component is always present regardless of the scan plane, i.e., it does not only occur when the scanner is looking at the specular angle. In addition to these components, we have also included a multiple-scattering approximation. It should be realized, however, that the multiple scattering applies only to an atmosphere with the sun as a source. Another multiple scattering calculation should be performed to include the effect due to the virtual sun.

The general result of all these calculations indicates that the various components are all about equal in magnitude but that there is considerable variation with respect to scan angle and visibility. Also, it appears that the multiply-scattered component is of major significance.

It must be pointed out that the objective of this investigation is to provide an algorithm for the correction of remotely sensed data for atmospheric effects so that one can extract from the multi-spectral data the radiance which is characteristic of the water itself. In our investigation we have dealt with the water surface as a flat, specular reflector, which in general is not true. A wind-roughened surface will be characterized by a complex wave structure which leads to a more complicated representation of the reflected and virtual sun radiances than presented in this report. A further investigation should be conducted to model the water surface in terms of wind speed and a stochastic representation of the reflecting facets of the water surface. If this is done, then a more realistic model could be developed which should provide better values for the sky-reflected and the virtual sun path radiances. It may even be possible to establish a method for the determination of wind speed by observing the average radiance as a function of the instantaneous field of view.

A further recommendation is to improve the accuracy of the algorithm by including a more detailed calculation of the multiply-scattered path radiance, both for the direct sun as a source and for the virtual sun as a source.

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